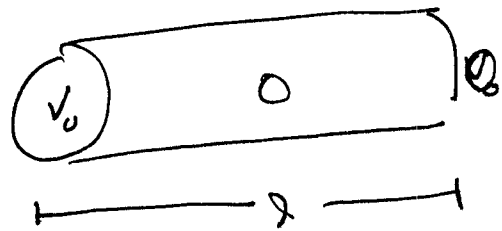
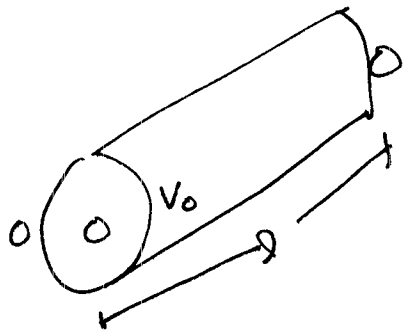


Bessel Functions

Let's relax the restriction that z is uniform.

Consider the systems



Laplace Eqn Cylindrical Coordinates

$$\frac{\partial^2 V}{\partial \rho^2} + \frac{1}{\rho} \frac{\partial V}{\partial \rho} + \frac{1}{\rho^2} \frac{\partial^2 V}{\partial \phi^2} + \frac{\partial^2 V}{\partial z^2} = 0$$

Separated Solution

$$V(\rho, \phi, z) = P(\rho) \Phi(\phi) Z(z) = R(\rho, \phi) Z(z)$$

Separate $Z(z)$ first to give

(2)

$$\frac{1}{R} \left(\frac{\partial^2 R}{\partial \rho^2} + \frac{1}{\rho} \frac{\partial R}{\partial \rho} + \frac{1}{\rho^2} \frac{\partial^2 R}{\partial \phi^2} \right) + \frac{1}{Z} \frac{\partial^2 Z}{\partial z^2} = 0$$

$$\frac{d^2 Z}{dz^2} - k^2 Z = 0$$

$$Z = \cancel{e^{ikz} \text{ or } \sin kz, \cos kz} \\ e^{\pm kz}$$

Now we have

$$\frac{\partial^2 R}{\partial \rho^2} + \frac{1}{\rho} \frac{\partial R}{\partial \rho} + \frac{1}{\rho^2} \frac{\partial^2 R}{\partial \phi^2} + k^2 R = 0$$

$$\rho^2 \frac{\partial^2 R}{\partial \rho^2} + \rho \frac{\partial R}{\partial \rho} + \rho^2 k^2 R + \frac{\partial^2 R}{\partial \phi^2} = 0$$

Separate Again

$$R = P(\rho) \Phi(\phi)$$

(3)

$$\frac{1}{P} \left[p^2 \frac{\partial^2 P}{\partial p^2} + p \frac{\partial P}{\partial p} + p^2 k^2 P \right] + \frac{1}{\Phi} \frac{\partial^2 \Phi}{\partial \phi^2} = 0$$

+ ν^2
- ν^2

where ν integer

$$\frac{d^2 \Phi}{d\phi^2} + \nu^2 \Phi = 0 \quad \Phi = \begin{matrix} \sin \nu \phi \\ \cos \nu \phi \end{matrix}$$

Leaving

$$p^2 \frac{\partial^2 P}{\partial p^2} + p \frac{\partial P}{\partial p} + p^2 k^2 P = \nu^2 P$$

$$\frac{\partial^2 P}{\partial p^2} + \frac{1}{p} \frac{\partial P}{\partial p} + \left(k^2 - \frac{\nu^2}{p^2} \right) P = 0$$

$$\text{Let } x = \frac{x_{0n} \rho}{a} \quad dx = \frac{x_{0n}}{a} d\rho$$

$$\rho = \frac{0}{x_{0n}}$$

$$A_n = \frac{2V_0}{a^2 J_1^2(x_{0n}) \sinh(k_{0n} l)} \left(\frac{a}{x_{0n}}\right)^2 \int_0^{x_{0n}} x J_0(x) dx$$

$$\int_0^{x_{0n}} x J_0(x) dx = x J_1(x) \Big|_0^{x_{0n}} = x_{0n} J_1(x_{0n})$$

$$A_n = \frac{2V_0}{x_{0n} J_1(x_{0n}) \sinh(k_{0n} l)}$$

Let $x = kp$

$$\frac{d^2 P}{dx^2} + \frac{1}{x} \frac{dP}{dx} + \left(1 - \frac{\nu^2}{x^2}\right) P = 0$$

\Rightarrow Bessel's Eyn

Solutions Bessel Functions

~~$J_{\nu}(x)$~~ $J_{\nu}(x)$

Neumann Functions

$\#$ $N_{\nu}(x)$

Gamma Function

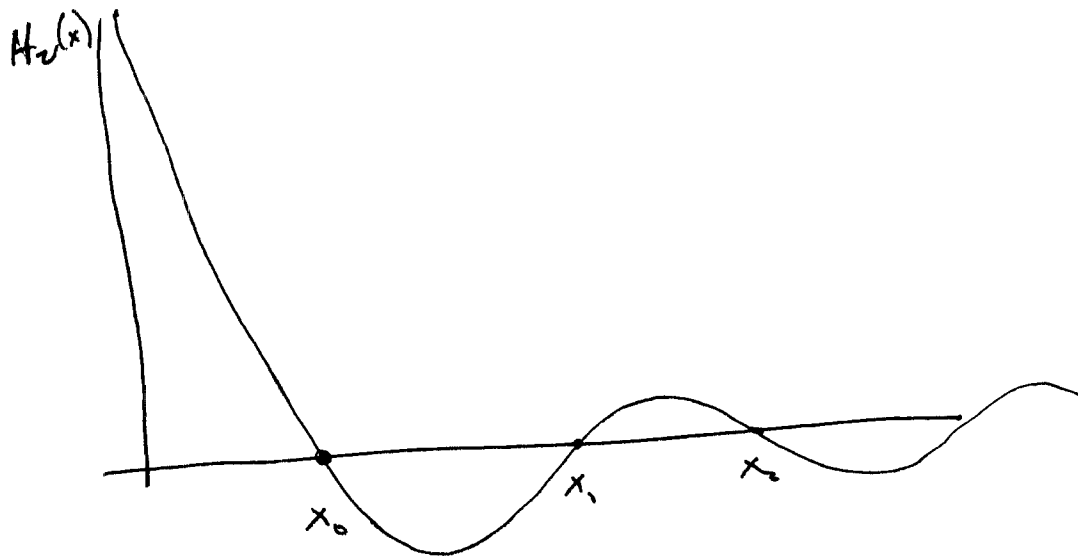
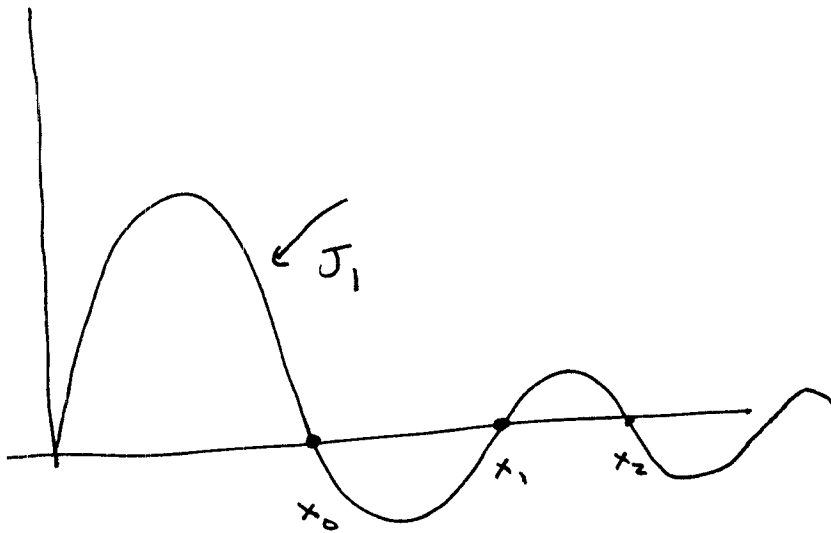
$$\Gamma(x) = (x-1)!$$

Limits

$$J_{\nu}(x) = \begin{cases} \frac{1}{\Gamma(\nu+1)} \left(\frac{x}{2}\right)^{\nu} & x \ll 1 \\ \sqrt{\frac{2}{\pi x}} \cos\left(x - \frac{\nu\pi}{2} - \frac{\pi}{4}\right) & x \gg 1 \end{cases}$$

$$N_{\nu}(x) = \begin{cases} \frac{2}{\pi} \left[\ln\left(\frac{x}{2}\right) + 0.5772 + \dots \right] & \nu=0 \quad x \ll 1 \\ -\frac{\Gamma(\nu)}{\pi} \left(\frac{x}{2}\right)^{\nu} & \nu \neq 0 \quad x \ll 1 \\ \sqrt{\frac{2}{\pi x}} \sin\left(x - \frac{\nu\pi}{2} - \frac{\pi}{4}\right) & x \gg 1 \end{cases}$$

$J_\nu(x)$



Each Bessel function has a set of zero

$$\nu=0 \quad x_{01} = 2.405 \quad x_{02} = 5.520 \dots$$

$$\nu=1 \quad x_{11} = 3.832 \quad x_{12} = 7.016 \dots$$

$$J_\nu(x_{\nu n}) = 0$$

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Orthogonality

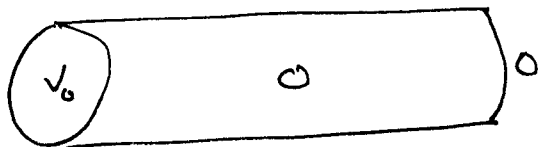
$$\int_0^a p J_\nu\left(x_{\nu n} \frac{p}{a}\right) J_\nu\left(x_{\nu m} \frac{p}{a}\right) dp \\ = \frac{a^2}{2} \left[J_{\nu+1}(x_{\nu n}) \right]^2 \delta_{mn}$$

Completeness

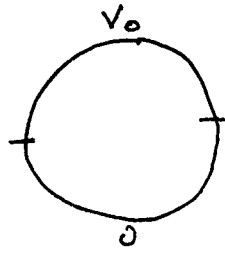
$$f(p) = \sum_{n=1}^{\infty} A_{\nu n} J_\nu\left(x_{\nu n} \frac{p}{a}\right)$$

$$A_{\nu n} = \frac{2}{a^2 J_{\nu+1}^2(x_{\nu n})} \int_0^a p f(p) J_\nu\left(\frac{x_{\nu n} p}{a}\right) dp$$

So these can be used to solve



The other case we considered is



of length l with both ends zero.

This time the Z separation must yield sines and cosines.

$$\frac{1}{Z} \frac{d^2 Z}{dz^2} = -k^2$$

Gives

$$\frac{d^2 R}{dx^2} + \frac{1}{x} \frac{dR}{dx} - \left(1 + \frac{\nu^2}{x^2}\right) R = 0$$

\Rightarrow Solutions Modified Bessel Functions

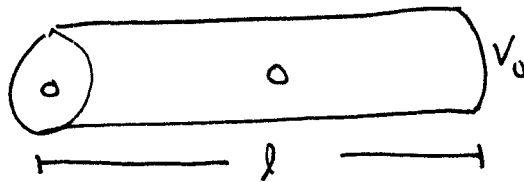
$$I_\nu(x) = i^{-\nu} J_\nu(ix)$$

$$K_\nu(x) = \frac{\pi}{2} i^{\nu+1} H_\nu^{(1)}(ix)$$

\nwarrow Hankel Function

$$H_\nu^{(1)}(x) = J_\nu(x) + iN_\nu(x)$$

For our cylinder,



our possible solutions are $l, \ln p, \phi, z$

$$[J_\nu(\rho k) + N_\nu(\rho k)] \times [\sin \nu \phi, \cos \nu \phi] [e^{kz}, e^{-kz}]$$

All the trivial solutions can be eliminated. N_ν blows up.

To satisfy $V=0$ at $z=0$ take $\frac{e^{kz} - e^{-kz}}{z} = \sinh kz$

for the z dependence.

There is no ϕ dependence, so used $\nu=0$ solution. This

leaves us

$J_0(\rho k) \sinh kz$ for our solutions

Now satisfy the boundary condition $V(a, \phi) = 0$.

$$\Rightarrow J_0(ak) = 0$$

$$\Rightarrow k_{0n} = \frac{x_{0n}}{a} \quad \text{where } x_{mn} \text{ are the zeros of the Bessel functions.}$$

The complete solution is then

$$V(\rho, \phi, z) = \sum_n A_n J_0(x_{0n} \rho/a) \sinh(k_{0n} z)$$

Satisfy final boundary condition

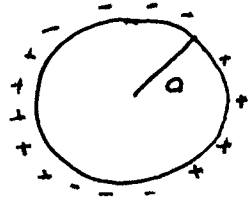
$$V(\rho, \phi, l) = V_0 = \sum_n A_n J_0(x_{0n} \rho/a) \sinh(k_{0n} l)$$

$$A_n = \frac{2V_0}{a^2 J_1^2(x_{0n}) \sinh(k_{0n} l)} \int_0^a \rho J_0\left(\frac{x_{0n} \rho}{a}\right) d\rho$$

$$A_n = \frac{2V_0}{a^2 J_1^2(x_{0n}) \sinh(k_{0n} l)}$$

E_x An infinitely long cylinder has a potential at the surface of $V(a, \phi) = V_0 \cos 2\phi$

Compute field everywhere.



Solutions $1, \ln \rho, \phi$

$$\rho^n \cos n\phi$$

$$\rho^n \sin n\phi$$

$$\rho^{-n} \cos n\phi$$

$$\rho^{-n} \sin n\phi$$

Inside

$$V(\rho, \phi) = \sum_n A_n \rho^n \cos n\phi + B_n \rho^n \sin n\phi$$

$$V(a, \phi) = V_0 \cos 2\phi = \sum A_n a^n \cos n\phi + B_n a^n \sin n\phi$$

$$\Rightarrow A_n, B_n = 0 \text{ except } A_2 = \frac{V_0}{a^2}$$

$$V(\rho, \phi)_{\text{inside}} = \frac{V_0}{a^2} \rho^2 \cos 2\phi$$

Outside

$$V(\rho, \phi) = \sum A_n \rho^{-n} \cos n\phi + B_n \rho^{-n} \sin n\phi$$

$$V(a, \phi) = V_0 \cos 2\phi$$

$$A_n, B_n = 0 \text{ except}$$

$$A_2 = a^2 V_0$$

$$V(\rho, \phi) = \frac{V_0 a^2}{\rho^2} \cos 2\phi$$

Field Inside

$$\nabla f = \frac{\partial f}{\partial \rho} \hat{\rho} + \frac{1}{\rho} \frac{\partial f}{\partial \phi} \hat{\phi}$$

$$\vec{E}_{\text{inside}} = -\nabla V$$

$$= -\frac{2V_0}{a^2} \rho \cos 2\phi \hat{\rho} + \frac{2V_0}{a^2} \rho \sin 2\phi \hat{\phi}$$

$$\vec{E}_{\text{outside}} = \frac{2V_0 a^2}{\rho^3} \cos 2\phi \hat{\rho} + \frac{2V_0 a^2}{\rho^3} \sin 2\phi \hat{\phi}$$

Surface Charge Use Pill box at surface

$$\sigma = (\vec{E}_{\text{outside}} - \vec{E}_{\text{inside}}) \cdot \hat{\rho} \epsilon_0$$

$$\sigma = \frac{2V_0}{a} \cos 2\phi - \left(-\frac{2V_0}{a} \cos 2\phi \right)$$

$$= \frac{4V_0}{a} \cos 2\phi$$