

## Lecture 6

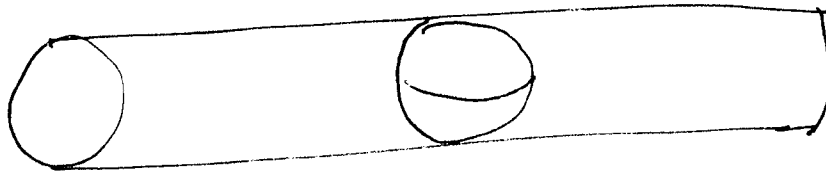
Linear Superposition and Odds and Ends.

Law of Linear Superposition - The total field at point  $P$  is the sum of the fields of the individual charges at  $P$

$$\vec{E}_P = \sum_i \vec{E}_{iP}$$

---

Consider a long tube of uniform charge density  $\rho$



and radius  $a$   
with a spherical  
cavity of radius  $a$ .

Compute field everywhere.

$$\vec{E} = \int \frac{\rho}{4\pi\epsilon_0} \frac{\hat{r}''}{r''^2} dv$$

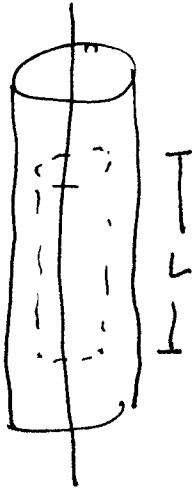
Ouch!

(2)

View system as a superposition of a tube with  $+p$  and a spherical volume charge with  $-p$ .

$$\vec{E}(\vec{r}) = \vec{E}_{\text{tube}} + \vec{E}_{\text{sphere}}$$

Pick center of the sphere as the origin and the axis of the tube as the  $z$ -axis.



Field of the Tube

Apply Gauss' law to a co-axial cylindrical of radius  $r$  and length  $L$

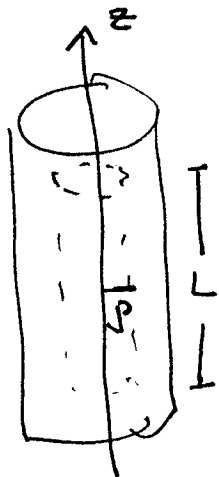
$$\text{If } r < 0, \quad Q_{\text{enc}} = 2\pi r L p$$

$$\Phi = 2\pi r L E = \frac{Q_{\text{enc}}}{\epsilon_0} = 2\pi r^2 L p$$

$$\vec{E} = \frac{pr}{2\epsilon_0} \hat{z}$$

Using  $\rho$  as the volume charge density here 3  
 introduces severe notational problems, let  $\rho \rightarrow \gamma$   
 and reserve  $\rho$  for cylindrical coordinate.

Field of Tube - Apply Gauss' law to a co-axial  
 cylindrical Gaussian surface of radius  $\rho$  and  
 length  $L$ .



$$\Phi = 2\pi\rho L E = \frac{Q_{enc}}{\epsilon_0} \Rightarrow E(\rho) = \frac{Q_{enc}}{2\pi\rho L \epsilon_0}$$

If  $\rho < a$ ,  $Q_{enc} = \pi\rho^2 L \gamma$

$$\vec{E}(\rho) = \frac{\pi\rho^2\gamma}{2\pi\rho\epsilon_0} \hat{\rho} = \frac{\gamma\rho}{2\epsilon_0} \hat{\rho}$$

If  $\rho \geq a$ ,  $Q_{enc} = \pi a^2 L \gamma$

$$\vec{E}(\rho) = \frac{\pi a^2 \gamma}{2\pi\epsilon_0\rho} \hat{\rho} = \frac{a^2\gamma}{2\epsilon_0\rho} \hat{\rho}$$

④

Likewise for the sphere,

$$\vec{E} = -\frac{r\gamma}{3\epsilon_0} \hat{r} \quad r < a$$

$$\vec{E} = -\frac{\frac{4}{3}\pi a^3 \gamma}{4\pi\epsilon_0 r^2} \hat{r} = -\frac{a^3 \gamma}{3\epsilon_0 r^2} \hat{r} \quad r > a$$

From these four fields we can construct the total field

$$p \geq a \quad \vec{E} = \frac{a^2 \gamma}{2\epsilon_0 p} \hat{p} \rightarrow \frac{a^3 \gamma}{3\epsilon_0 r^2} \hat{r}$$

$$r < a \quad \vec{E} = \frac{\gamma p}{2\epsilon_0} \hat{p} \rightarrow \frac{r\gamma}{3\epsilon_0} \hat{r}$$

$$r > a, p < a \quad \vec{E} = -\frac{a^3 \gamma}{3\epsilon_0 r^2} \hat{r} \rightarrow \frac{\gamma p}{2\epsilon_0} \hat{p}$$

# Lecture 6 - Dipoles

14

So we can do anything

## Linear Charges ( $\lambda$ )

$$\vec{E} = \int \frac{\lambda d\vec{l} \hat{r}''}{4\pi\epsilon_0 r''^2}$$

$$V = \int \frac{\lambda dl}{4\pi\epsilon_0 r''}$$

## Surface Charges ( $\sigma$ )

$$\vec{E} = \int \frac{\sigma da \hat{r}''}{4\pi\epsilon_0 r''^2}$$

$$V = \int \frac{\sigma da}{4\pi\epsilon_0 r''}$$

## Volume Charges ( $\rho$ )

$$\vec{E} = \int \frac{\rho dv \hat{r}''}{4\pi\epsilon_0 r''^2}$$

$$V = \int \frac{\rho dv}{4\pi\epsilon_0 r''}$$

## Green's Functions

The potential presents us with a fairly general problem. We have a set of "sources"  $p(x)$  and a differential equation  $Lu(x) = p(x)$  that defines how the sources act. Our job is to find  $u(x)$ .

It would be nice to write  $u(x)$  as a sum over the sources

$$u(x) = \int G(x,s) p(s) ds$$

where  $G(x,s)$  is some function that "propagates" the effect of a source at  $s$  to the point  $x$ .

The function  $u(x)$  must still solve the equation

$$\text{so } Lu(x) = L \int G(x,s) p(s) ds$$

$$= \int [L G(x,s)] p(s) ds$$

because  $L$  acts on  $x$  not  $s$ .

$$\text{evidently } L G(x,s) = \delta(x-s)$$

(2)

$G(x, s)$  is called the Green's function of the operator  $L$  (the propagator in particle physics. The Green's functions are the lines in the Feynman diagrams).

Our problem is to solve  $\nabla \cdot \vec{E} = \rho/\epsilon_0$   
 $\nabla \times \vec{E} = 0$

We can solve the second equation if  $\vec{E} = -\nabla V$  where  $V$  is some function. This turns the first equation into  $\nabla \cdot (-\nabla V) = \rho/\epsilon_0$  or

$$\nabla^2 V = -\rho/\epsilon_0 \quad (\text{Poisson's Eqn})$$

Solve  $-\epsilon_0 \nabla^2 V = \rho \Rightarrow$  We want the Green's function of  $L = -\epsilon_0 \nabla^2$

Let's drop the  $\epsilon_0$ , since the Green's function for the Laplacian is a pretty general problem.

$$-\nabla^2 (G(\vec{r}, \vec{r}')) = \delta(\vec{r} - \vec{r}')$$

Note  $\nabla$  acts on  $x, y, z$  not  $x', y', z'$

Looking back, we have

$$-\nabla^2 \left( \frac{1}{r''} \right) = 4\pi \delta(\vec{r} - \vec{r}'')$$

$$G(\vec{r}, \vec{r}') = -\frac{1}{4\pi r''} = \frac{1}{4\pi r''}$$

$$r'' = \sqrt{(x-x')^2 + (y-y')^2 + (z-z')^2}$$



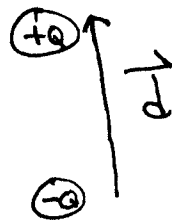
# Dipole Moment

There is a further method of specifying the charge of a system, by the dipole moment per unit volume.

Dipole System - Zero charge but non-zero separation of charge.

Dipole Moment

$$\vec{P} = Q \vec{d}$$



Polarization Vector ( $\vec{P}$ ) - Dipole moment per unit volume.

$$\text{Dipole Potential } (V_{dip}) = \frac{d \left( V(z - d/2) - V(z + d/2) \right)}{d}$$

$$= \frac{qd}{4\pi\epsilon_0} \left( \frac{1}{(x^2 + y^2 + (z - d/2)^2)^{1/2}} - \dots \right)$$

$$V_{dip} = \frac{qd}{4\pi\epsilon_0} \left( -\frac{\partial}{\partial z} \left( \frac{1}{r} \right) \right) \quad \lim_{d \rightarrow 0}$$

$$qd = p$$

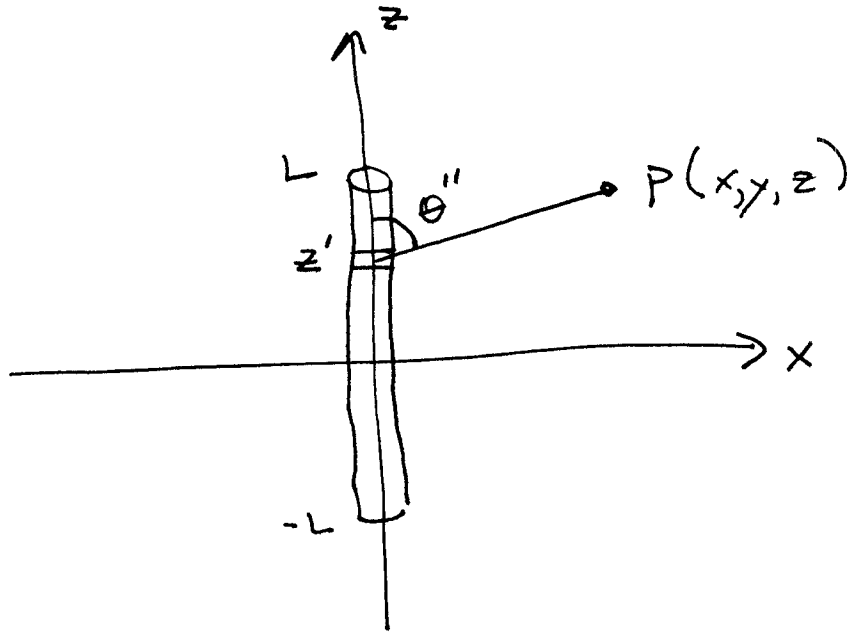
$$V_{dip} = \frac{p}{4\pi\epsilon_0} \left( \frac{z}{(\sqrt{x^2+y^2+z^2})^3} \right)$$

$$= \frac{p}{4\pi\epsilon_0} \frac{\cos\theta}{r^2} = \frac{k p \cos\theta}{r^2}$$

$$V_{dip} = \frac{\vec{p} \cdot \hat{r}}{4\pi\epsilon_0 r^2}$$

Calculation of field left as hwk problem.

Ex Compute the field of a thin pencil with polarization  $\vec{P} = P_0 \hat{z}$  and area  $A$  occupying  $-L$  to  $L$



$$\begin{aligned} \vec{r}'' &= \vec{r} - \vec{r}' = (x, y, z) - (0, 0, z') \\ &= (x, y, z - z') \end{aligned}$$

$$\cos \theta'' = \hat{z} \cdot \hat{r}'' = \frac{z - z'}{\sqrt{x^2 + y^2 + (z - z')^2}}$$

$$V = \int_{-L}^L \frac{k \phi p \cos \theta''}{(\sqrt{x^2 + y^2 + (z - z')^2})^2} dz' \quad dp = A P_0 dz'$$

$$= k A P_0 \int_{-L}^L \frac{(z - z') dz'}{(x^2 + y^2 + (z - z')^2)^{3/2}}$$

$$V_d = KAP_0 \left( \frac{1}{(x^2 + y^2 + (z-L)^2)^{1/2}} - \frac{1}{(x^2 + y^2 + (z+L)^2)^{1/2}} \right)$$

$$\begin{aligned}
 &> \int \left( \frac{(z-zp)}{(x^2+y^2+(z-zp)^2)^{\frac{3}{2}}}, zp \right); \\
 &> \frac{1}{\sqrt{x^2+y^2+z^2-2zzp+zp^2}}
 \end{aligned}
 \tag{1}$$

# Electric Force

$$\vec{F} = Q \vec{E}$$

$$\vec{F} = \int_C \vec{E} \lambda dl$$

$$\vec{F} = \int_S \vec{E} \sigma da$$

$$\vec{F} = \int_V \vec{E} \rho dv$$

Ex Force of  $q_1 = 3nC$  on  $q_2 = 5nC$

$$\vec{r}_1 = (3cm, 0, 0)$$

$$\vec{r}_2 = (0, 5cm, 0)$$

$$\vec{F}_{12} = \frac{kq_1 q_2}{r_{12}^2} \hat{r}_{12}$$

$$\vec{r}_{12} = \vec{r}_2 - \vec{r}_1 = (-3cm, 5cm, 0)$$

$$r_{12} = \sqrt{(3cm)^2 + (5cm)^2}$$

$$= \sqrt{34} cm$$

$$\hat{r}_{12} = \left( -\frac{3}{\sqrt{34}}, \frac{5}{\sqrt{34}}, 0 \right)$$

