

## Maxwell's Eqns for Magnetism

Let's compute  $\nabla \cdot \vec{B}$  and  $\nabla \times \vec{B}$  from the Biot-Savart Law.

$$\nabla \cdot \vec{B} = \frac{\mu_0}{4\pi} \int_V \nabla \cdot \left( \frac{\vec{J}(\vec{r}') \times \hat{r}''}{r''^2} \right) d\tau'$$

Identity

$$\nabla \cdot (\vec{A} \times \vec{B}) = \vec{B} \cdot (\nabla \times \vec{A}) - \vec{A} \cdot (\nabla \times \vec{B})$$

$$\begin{aligned} \nabla \cdot \left( \vec{J} \times \left( \frac{\hat{r}''}{r''^2} \right) \right) &= \frac{\hat{r}''}{r''^2} \cdot (\nabla \times \vec{J}) \\ &\quad - \vec{J} \cdot \left( \nabla \times \frac{\hat{r}''}{r''^2} \right) \end{aligned}$$

$\nabla \times \vec{J}(\vec{r}') = 0$  because  $\vec{J}$  depends only on  $\vec{r}'$  coordinates, but derivative taken with respect to un-primed coordinates

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$$\nabla \times \left( \frac{\hat{r}}{r^{1/2}} \right) = 0$$

because this is just the curl of the point charge radial dependence.

$$\text{so } \nabla \cdot \vec{B} = 0$$

### No Magnetic Monopoles

$$\bullet \nabla \cdot \vec{B} = 0 \iff \int \vec{B} \cdot d\vec{a} = \Phi_m = 0$$

• Total magnetic flux out of a closed surface is zero.

• No isolated magnetic charge.

• Magnetic field lines are closed curves.

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Examine  $\nabla \times \vec{B}$

$$\nabla \times \vec{B} = \frac{\mu_0}{4\pi} \int_V \nabla \times \left( \vec{J}(\vec{r}') \times \frac{\hat{r}''}{r''^2} \right) d\tau'$$

Identity

$$\begin{aligned} \nabla \times (\vec{A} \times \vec{B}) &= (\vec{B} \cdot \nabla) \vec{A} - (\vec{A} \cdot \nabla) \vec{B} \\ &\quad + \vec{A} (\nabla \cdot \vec{B}) - \vec{B} (\nabla \cdot \vec{A}) \end{aligned}$$

$$\begin{aligned} \nabla \times \vec{B} &= \frac{\mu_0}{4\pi} \int_V \left[ \overset{\textcircled{1}}{\left( \frac{\hat{r}''}{r''^2} \cdot \nabla \right) \vec{J}} - \overset{\textcircled{2}}{\left( \vec{J} \cdot \nabla \right) \left( \frac{\hat{r}''}{r''^2} \right)} \right. \\ &\quad \left. + \overset{\textcircled{3}}{\vec{J} (\nabla \cdot \left( \frac{\hat{r}''}{r''^2} \right))} - \overset{\textcircled{4}}{\frac{\hat{r}''}{r''^2} (\nabla \cdot \vec{J})} \right] d\tau' \end{aligned}$$

$\textcircled{1} + \textcircled{4} = 0$  because derivatives taken with respect to unprimed variables, but  $\vec{J}(\vec{r}')$  depends only on primed variables.

From before,

$$\nabla \cdot \left( \frac{\hat{r}''}{r''^2} \right) = 4\pi \delta(\vec{r}'')$$

so ③ =  $\vec{J}(\vec{r}') 4\pi \delta(\vec{r}'')$

Now work on piece ②

$$\int_V (\vec{J} \cdot \nabla) \left( \frac{\hat{r}''}{r''^2} \right) d\tau'$$

In  $r''$ ,  $\hat{r}''$  the variables  $x, x', y, y', z, z'$  only appear in combinations  $x-x', y-y', z-z'$   
 So the effect of the transformation  $x \rightarrow x'$  is to change signs.

$$(\vec{J} \cdot \nabla) \left( \frac{\hat{r}''}{r''^2} \right) = - (\vec{J} \cdot \nabla') \left( \frac{\hat{r}''}{r''^2} \right)$$

More identities,

$$\nabla \cdot (f \vec{A}) = f(\nabla \cdot \vec{A}) + (\vec{A} \cdot \nabla) f$$

$$(\vec{A} \cdot \nabla) f = \nabla \cdot (f \vec{A}) - f(\nabla \cdot \vec{A})$$

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This is relatively complicated, let's look at the x-component. The other components will be similar.

$$\begin{aligned} (\vec{J}' \cdot \nabla') \left( \frac{x-x'}{r'^3} \right) &= \nabla' \cdot \left( \frac{x-x'}{r'^3} \cdot \vec{J}'(\vec{r}') \right) \\ &\quad - \frac{x-x'}{r'^3} (\nabla' \cdot \vec{J}'(\vec{r}')) \end{aligned}$$

The second term is zero for magnetostatics,  $\nabla' \cdot \vec{J}' = 0$ .

Integrate the remaining part of ②

$$\int_V \nabla' \cdot \left( \frac{x-x'}{r'^3} \vec{J}'(\vec{r}') \right) d\vec{r}'$$

$$= \int_S \frac{x-x'}{r'^3} \vec{J}'(\vec{r}') d\vec{r}'$$

$\Rightarrow 0$  if we ~~ensure~~ ensure  $S$  encloses all charge, and so no current passes through the surface.

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So all that is left is ③

$$\begin{aligned}\nabla \times \vec{B} &= \frac{\mu_0}{4\pi} \int_V \vec{J}(\vec{r}') 4\pi \sigma(\vec{r}'') d\tau' \\ &= \mu_0 \vec{J}(\vec{r})\end{aligned}$$

Ampere's Law (magnetostatics)

$$\nabla \times \vec{B} = \mu_0 \vec{J}$$

Convert to integral form

$$\int_S (\nabla \times \vec{B}) \cdot d\vec{a} = \int_S \mu_0 \vec{J} \cdot d\vec{a}$$

$$d\vec{a} = \hat{n} da$$

Stokes's Thm

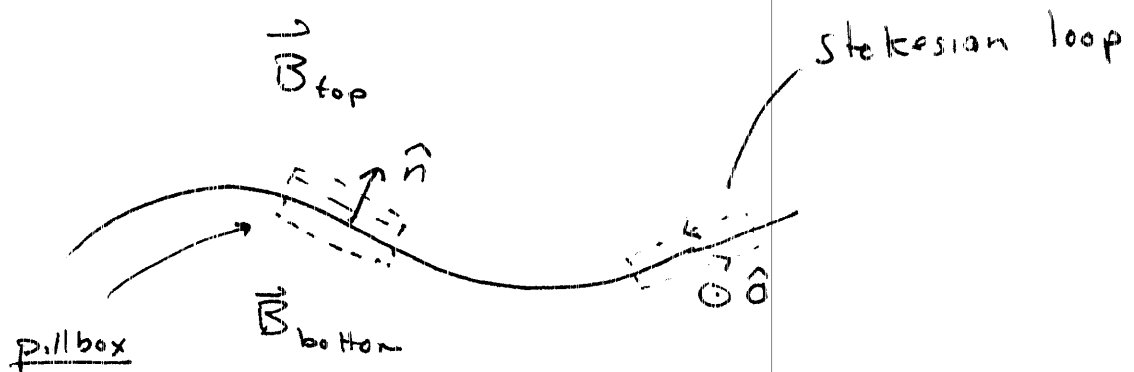
$$\int_C \vec{B} \cdot d\vec{\ell} = \mu_0 I_{enc}$$

Ampere's Law - Integral Form

$$\int_C \vec{B} \cdot d\vec{\ell} = \mu_0 I_{enc}$$

# Magnetostatic Boundary Conditions

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Apply  $\nabla \cdot \vec{B}$  to pillbox

$$\Phi_m = \vec{B}_{top} \cdot \hat{n} A - \vec{B}_{bottom} \cdot \hat{n} A = 0$$

$$B_{top}^{\perp} = B_{bottom}^{\perp}$$

Apply  $\nabla \times \vec{B} = \mu_0 \vec{J}_{enc}$  to Stokesian Loop

$$\vec{B}_{top} \cdot \hat{t} l - \vec{B}_{bottom} \cdot \hat{t} l = \mu_0 \vec{K} \cdot \hat{a}$$

where  $\vec{K}$  is the surface current and  $\hat{a}$  is the surface normal to the curve bounded by the Stokesian loop.

In vector form these can be combined as

$$\vec{B}_{top} - \vec{B}_{bottom} = \mu_0 (\vec{K} \times \hat{n})$$