

## Review Lecture

Ex Estimate the classical radius of an electron by setting its relativistic energy

$U = mc^2$ ,  $m_e = 9.11 \times 10^{-31} \text{ kg}$ ,  $c = 3.00 \times 10^8 \text{ m/s}$   
equal to electrostatic energy stored in its field  
under the assumption that the electron's charge  
 $q_e = -1.602 \times 10^{-19} \text{ C}$  is on the surface of a  
sphere of radius  $a$ .

Sol Find energy by integrating energy density

$$U = \int_{\text{space}} \frac{1}{2} \epsilon_0 E^2 d\tau$$

$$E = \frac{k q_e}{r^2}$$

$$U = \left( \frac{1}{2} \epsilon_0 \right) \left( \frac{1}{4\pi\epsilon_0} \right)^2 q_e^2 \int_a^\infty dr \int_0^{2\pi} d\phi \int_0^\pi d\theta r^2 \sin\theta \left( \frac{1}{r^2} \right)^2$$

$$= \frac{q_e^2}{32 \epsilon_0 \pi^2} \cdot 4\pi \int_a^\infty \frac{dr}{r^2}$$

$$= \frac{q_e^2}{8 \epsilon_0 \pi} \left( -\frac{1}{r} \right)_a^\infty = \frac{q_e^2}{8 \epsilon_0 \pi a}$$

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Check using capacitance

$$C = 4\pi\epsilon_0 d$$

$$U = \frac{1}{2} \frac{q_e^2}{C} = \frac{q_e^2}{8\pi\epsilon_0 d}$$

$$mc^2 = \frac{q_e^2}{8\pi\epsilon_0 d}$$

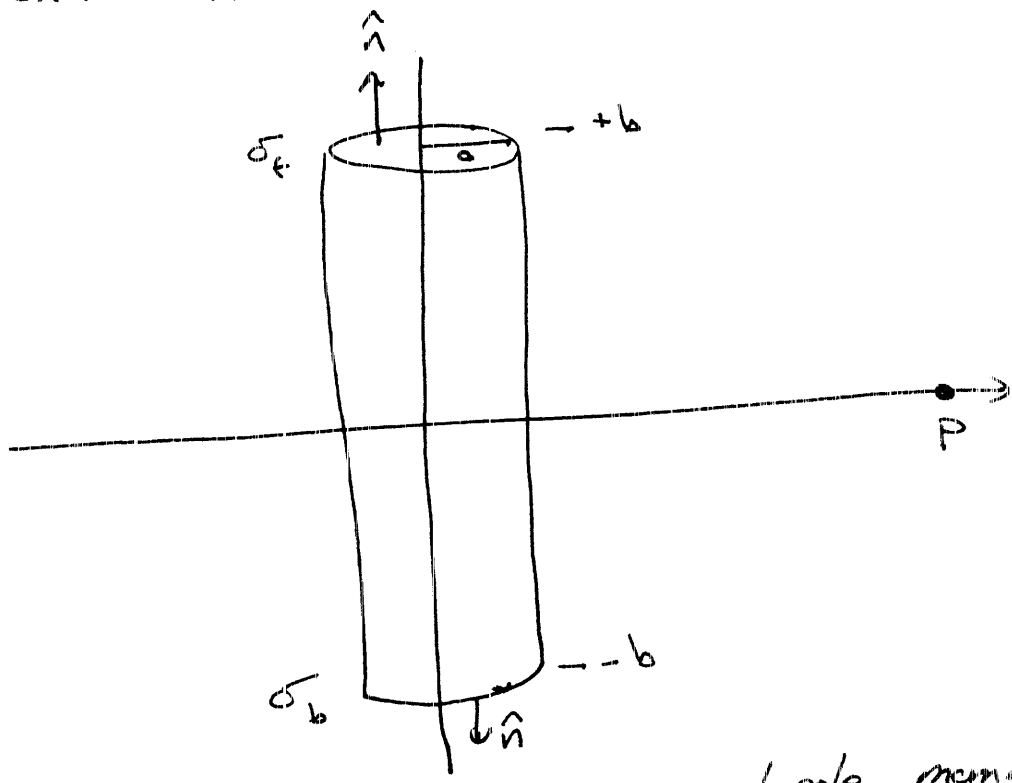
$$d = \frac{q_e^2}{8\pi\epsilon_0 mc^2} = 1.4 \times 10^{-15} \text{ m}$$

(about the size of a proton).

③

E<sub>x</sub> A cylinder of radius  $a$  and length  $2b$  is centered on the origin. The axis of the cylinder is the  $z$ -axis. The cylinder contains a material with polarization  $\vec{P} = \frac{\gamma}{2} z^2 \hat{z}$

Compute the field at a point  $10b$  along the  $x$ -axis (approximately).  $b \gg a$ .



Find approximate field using dipole moment.

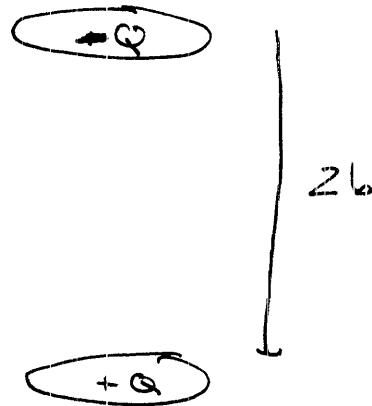
$$\rho_b = -\nabla \cdot \vec{P} = \gamma z$$

$$\sigma_b = \vec{P} \cdot (-\hat{z}) = \frac{1}{2} \gamma b^2$$

$$\sigma_t = \vec{P} \cdot (\hat{z}) = -\frac{1}{2} \gamma b^2$$

The top and bottom form a simple dipole

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$$Q = \sigma_b \pi a^2 = \frac{1}{2} \gamma b^2 \cdot \pi a^2 = \frac{1}{2} \pi \gamma a^2 b^2$$

$$\vec{p} = -Q d \hat{z} = -\left(\frac{1}{2} \pi \gamma a^2 b^2\right)(2b) \hat{z}$$

$$= -\pi \gamma a^2 b^3 \hat{z}$$

The dipole moment of the volume bound charge

is

$$\vec{p} = \int_{\text{cylinder}} \vec{r}' \rho(\vec{r}') d\tau'$$

The dipole moment evidently points in the  $\hat{z}$  direction, so work on the  $z$  component of  $p$ .

(5)

$$P_z = \int z' p_b(z') d\tau'$$

$$= \int z' (+\gamma z') d\tau'$$

$$= \gamma \int z'^2 d\tau' \quad d\tau' = dz' ds' s' d\phi'$$

$$= \gamma \int_{-b}^b dz' \int_0^{2\pi} d\phi' \int_0^a ds' s' z'^2$$

$$= 2\pi\gamma \cdot \frac{a^2}{2} \cdot \int_{-b}^b dz' z'^2$$

$$= \pi\gamma a^2 \left( \frac{z'^3}{3} \right)_{-b}^b = \frac{2}{3} \pi\gamma a^2 b^3$$

$$\vec{P}_{\text{bulk}} = \frac{2}{3} \pi\gamma a^2 b^3 \hat{z}$$

Total dipole moment

$$\vec{P} = \vec{P}_{\text{bulk}} + \vec{P}_{\text{surface}} = \frac{2}{3} \pi\gamma a^2 b^3 \hat{z} - \pi\gamma a^2 b^3 \hat{z}$$

$$= -\frac{1}{3} \pi\gamma a^2 b^3 \hat{z}$$

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Field at a point on x-axis

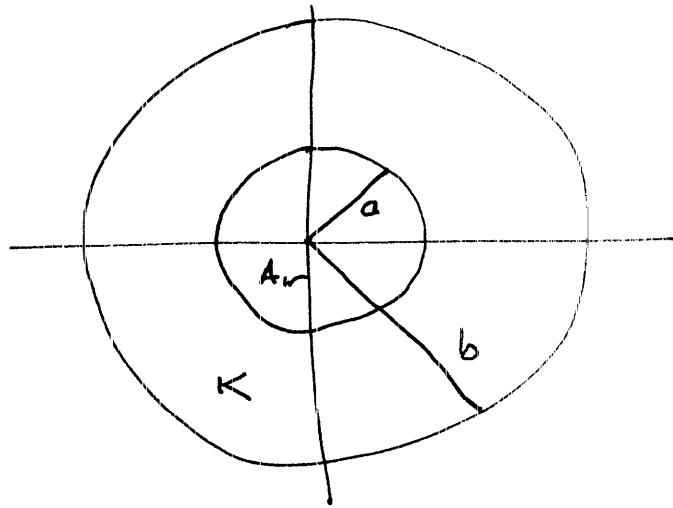
$$\vec{E} = -\frac{p k}{x^3} \hat{z}$$

$$= -\frac{p k}{(10b)^3} \hat{z} = -\frac{1}{4\pi\epsilon_0} \cdot \frac{1}{1000b^3} \cdot \left(-\frac{1}{3} \pi \gamma a^2 b^3\right)$$

$$= \frac{1}{12,000} \cdot \frac{\gamma a^2}{\epsilon_0} \hat{z}$$

(7)

Ex The surface of a thick dielectric cylinder is held at  $V(b, \phi) = V_0 \sin 2\phi$ . The inner radius of the shell is  $a$  and the shell has dielectric constant  $\kappa = \epsilon_r$ . Find the bound charge distribution on the inner surface of the shell.



Solve Laplace's Eqn

For  $s < a$

$$V_{\text{air}} = \sum_n A_n \sin n\phi s^n$$

For  $a < s < b$

$$V_{\kappa} = \sum_n B_n \sin n\phi s^n + C_n \sin n\phi s^{-n}$$

where I've thrown away  $\cos n\phi$  solutions because BC depends only on  $\sin n\phi$ .

Continuity BC

$$V_{\kappa}(b, \phi) = V_0 \sin 2\phi$$

$$V_{air}(a, \phi) = V_{\kappa}(a, \phi)$$

ES BC

$$\epsilon_{\kappa} \frac{\partial V_{\kappa}}{\partial s} \Big|_a - \epsilon_{air} \frac{\partial V_{air}}{\partial s} \Big|_a = -\sigma_f = 0$$

$$\epsilon_{\kappa} = \kappa \epsilon_0$$

$$\epsilon_{air} = \epsilon_0$$

$$\kappa \frac{\partial V_{\kappa}}{\partial s} \Big|_a - \frac{\partial V_{air}}{\partial s} \Big|_a = 0$$

Apply BC

$$V_{\kappa}(b, \phi) = V_0 \sin 2\phi = \sum B_n \sin n\phi b^n + C_n \sin n\phi b^{-n}$$

$$\Rightarrow B, C = 0 \quad n \neq 2$$

$$V_0 = B_2 b^2 + C_2 b^{-2}$$



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$$V_{\text{air}}(a, \phi) = V_{\kappa}(a, \phi)$$

$$\sum A_n a^n \sin n\phi = B_2 a^2 \sin 2\phi + \frac{C_2}{a^2} \sin 2\phi$$

Again orthogonality implies,  $A_n = 0 \quad n \neq 2$

$$A_2 a^2 = B_2 a^2 + \frac{C_2}{a^2}$$

$$A_2 = B_2 + \frac{C_2}{a^4}$$

Electrostatic BC	$V_{\text{air}} = A_2 s^2 \sin 2\phi$ $V_{\kappa} = B_2 s^2 \sin 2\phi + C_2 s^{-2} \sin 2\phi$
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$$\left. \frac{\partial V_{\kappa}}{\partial s} \right|_a = 2B_2 a \sin 2\phi - \frac{2C_2}{a^3} \sin 2\phi$$

$$\left. \frac{\partial V_{\text{air}}}{\partial s} \right|_a = 2A_2 a \sin 2\phi$$

$$\kappa \left. \frac{\partial V_{\kappa}}{\partial s} \right|_a - \left. \frac{\partial V_{\text{air}}}{\partial s} \right|_a = \kappa \left( 2B_2 a \sin 2\phi - \frac{2C_2}{a^3} \sin 2\phi \right)$$

$$- 2A_2 a \sin 2\phi = 0$$

$$2B_2 a^k - \frac{2C_2 k}{a^3} - 2A_2 a = 0$$

$$B_2 k - \frac{C_2 k}{a^4} - A_2 = 0$$

Solve

$$B_2 + \frac{C_2}{b^4} = \frac{V_0}{b^2} \quad (1)$$

$$A_2 - B_2 - \frac{C_2}{a^4} = 0 \quad (2)$$

$$A_2 - B_2 k + \frac{C_2 k}{a^4} = 0 \quad (3)$$

$$(1) + (2) \Rightarrow A_2 = \frac{V_0}{b^2}$$

$$k(2) + (3) \Rightarrow$$

$$(k+1)A_2 - 2kB_2 = 0$$

$$B_2 = \frac{k+1}{2k} A_2 = \frac{k+1}{2k} \frac{V_0}{b^2}$$

$$C_2 = \sigma_A^4 (A_2 - B_2)$$

$$= \sigma^4 \left( \frac{V_0}{b^2} - \frac{\kappa+1}{2\kappa} \frac{V_0}{b^2} \right)$$

$$= \frac{V_0 \sigma^4}{b^2} \left( 1 - \frac{\kappa+1}{2\kappa} \right) =$$

$$= \frac{V_0 \sigma^4}{b^2} \left( \frac{2\kappa - \kappa - 1}{2\kappa} \right)$$

$$= \frac{V_0 \sigma^4}{b^2} \frac{\kappa - 1}{2\kappa}$$

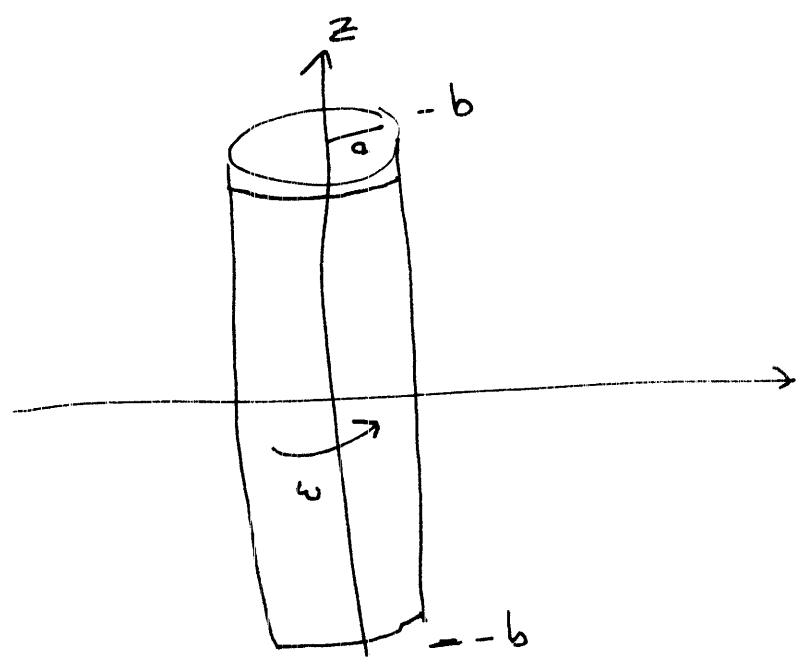
Charge Density at Surface

Use Pill box,

$$\Phi_e = (\vec{E}_\kappa \cdot \hat{s} + \vec{E}_{air} \cdot (-\hat{s})) A = \frac{\sigma A}{\epsilon_0}$$

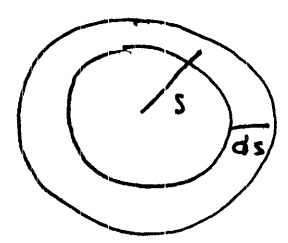
$$\left. \frac{\partial V_\kappa}{\partial s} \right|_a - \left. \frac{\partial V_{air}}{\partial s} \right|_a = \frac{\sigma}{\epsilon_0}$$

Ex A cylinder with charge density  $\rho = \gamma z$  rotates with angular velocity  $\omega$  about its axis. Compute field at a point far from the cylinder along  $z$ -axis.



Compute Moment of Ring

$$\vec{J} = \rho \vec{v} = \rho \omega s \hat{\phi}$$



$$d\vec{I} = \vec{J} \cdot d\vec{a} = \vec{J} ds ds$$

The moment of the ring is

$$dm = dIA = \pi s^2 dI = \pi s^2 J \omega s dz$$

The total moment is

$$m = \int dm = \int_{-b}^b dz \int_0^a ds (\pi s^2) (\omega s)$$

~~...~~

$$m = \pi \omega \int_{-b}^b dz \int_0^a ds s^2 \cdot s \cdot z$$

$$= \pi \omega \gamma \int_{-b}^b dz z \int_0^a ds s^3 = 0$$

~~...~~

Change problem so cylinder only extends from 0 to b.

$$m = \pi \omega \gamma \underbrace{\int_0^b dz z}_{\frac{b^2}{2}} \int_0^a ds s^3 = \frac{\pi \omega \gamma b^2 a^4}{8}$$

Field at a Point Along Axis

(14)

$$E = \frac{2\mu_0}{4\pi} \frac{m}{z^3} \hat{z}$$

Ex ~~Is~~  $\vec{B} = \gamma z^2 \hat{z}$  a valid  
 electrostatic magnetic field? If so find  
 the current that produced the field?

Sln  $\nabla \cdot \vec{B} = 2\gamma z \neq 0$  not ~~valid~~  
 valid field.

What about  $\vec{B} = \gamma z^2 \hat{\phi}$ ?

$$\nabla \cdot \vec{B} = 0 \quad \checkmark$$

$$\nabla \times \vec{B} = -\frac{\partial B_{\phi}}{\partial z} \hat{s} + \frac{1}{s} \frac{\partial (B_{\phi} s)}{\partial s} \hat{z}$$

$$= -\gamma z \hat{s} + \frac{\gamma z^2}{s} \hat{z} = \mu_0 \vec{J}$$

$$\vec{J} = \frac{1}{\mu_0} \left[ -\gamma z \hat{s} + \frac{\gamma z^2}{s} \hat{z} \right]$$

Check static current

$$\nabla \cdot \vec{J} = \frac{1}{\mu_0} \left[ \frac{1}{s} \frac{\partial}{\partial s} s J_s + \frac{\partial J_z}{\partial z} \right]$$

$$\nabla \cdot \vec{J} = \frac{1}{\mu_0} \left[ -\frac{\partial z}{\partial t} + \frac{z \partial z}{\partial t} \right] \neq 0$$

not static