

## Electromotive Force

Consider the work done by the electromagnetic fields as a charged particle moves along the path  $C$

$$\begin{aligned} \text{Work} &= \int_C (\vec{F}_e + \vec{F}_m) \cdot d\vec{l} \\ &= q \int_C (\vec{E} + \vec{v} \times \vec{B}) \cdot d\vec{l} \end{aligned}$$

Electromotive Force (emf) - Work per

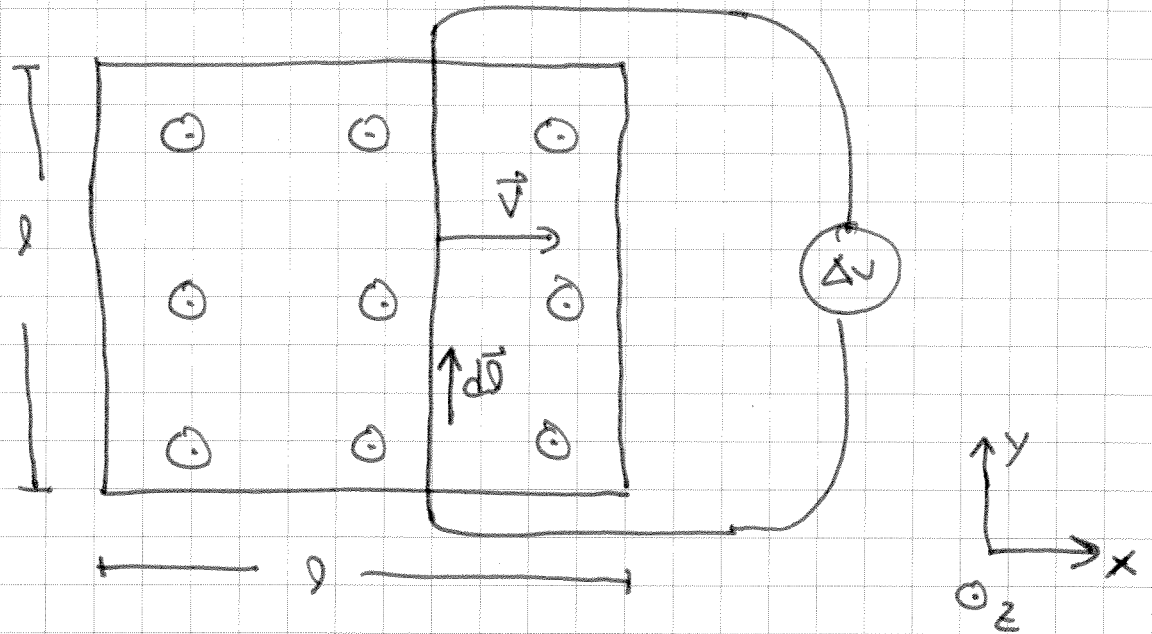
unit charge done by the electromagnetic fields as a charged particle moves along the path  $C$ .

$$\text{emf} = \frac{W}{q} = \int_C (\vec{E} + \vec{v} \times \vec{B}) \cdot d\vec{l}$$

Motional EMF The part of the emf resulting from the motion of the current path.

$$\text{emf} = \int_C (\vec{v} \times \vec{B}) \cdot d\vec{l}$$

Ex Pull wire through magnetic field



$$\vec{B} = B_0 \hat{z} \quad \vec{v} = v_0 \hat{x} \quad d\vec{l} = dy \hat{y}$$

$$\vec{v} \times \vec{B} = v_0 B_0 (\hat{x} \times \hat{z}) = -\hat{y} v_0 B_0$$

$$(\vec{v} \times \vec{B}) \cdot d\vec{l} = -v_0 B_0 \hat{y} \cdot \hat{y} dy = -v_0 B_0 dy$$

Motional EMF

$$\text{emf} = \int_C (\vec{v} \times \vec{B}) \cdot d\vec{l} = -v_0 B_0 l$$

⇒ We would measure this in Volts on a voltmeter.

Note, the directions are crucial, the magnetic force must be along the wire.

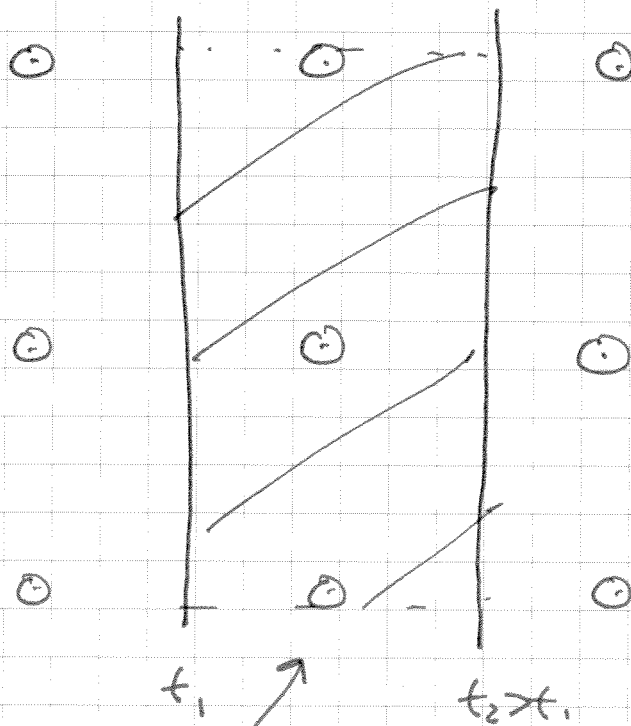
(1) If  $\vec{v}$  out of page,  $\vec{v} \times \vec{B} = 0$ ,  $\text{emf} = 0$ .

(2) If  $\vec{v}$  toward the top of the page,

$$\vec{v} = v_0 \hat{y} \Rightarrow \vec{v} \times \vec{B} = v_0 B_0 \hat{x}$$

$$\Rightarrow (\vec{v} \times \vec{B}) \cdot d\vec{\ell} = 0.$$

$\Rightarrow$  The key to having a non-zero emf is for the surface formed by the trajectory to have a non-zero flux through the surface.



surface traced by trajectory

Magnetic Flux  $\Phi_m$  The magnetic flux through surface  $S$  is

$$\Phi_m = \int_S \vec{B} \cdot d\vec{a}$$

⇒ If field is uniform and surface is flat,

$$\Phi_m = (\vec{B} \cdot \hat{n}) A = BA \cos \theta$$

↙ area of surface

$\theta$  - angle between  $\vec{B}$  and  $\hat{n}$ .

For our wire moving in the field at constant velocity, the flux traced out by the surface from time  $t_0$  to time  $t$  is

$$\Phi_m = B_0 l v_0 (t - t_0)$$

The emf is then the time rate of change of the flux

$$\text{emf} = - \frac{d\Phi_m}{dt} = -B_0 l v_0$$

Ex How many volts? Suppose  $B_0 = \frac{1}{4} \text{ T}$

$$v_0 = 10 \text{ m/s}, \quad \rho = 10 \text{ cm}$$

$$\text{emf} = B_0 v_0 \rho = \frac{1}{4} \text{ V}$$

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Multiple Turns ( $N$ ) - We will often deal with flux through surfaces formed by multiple turns of wire, in this case the flux becomes

$$\Phi_m = N \int_S \vec{B} \cdot d\vec{a}$$

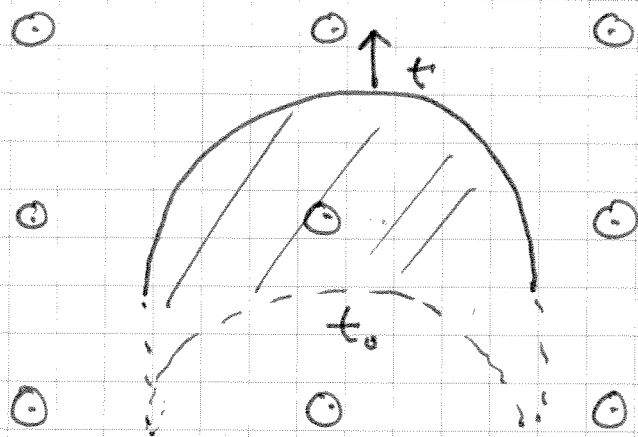
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Flux Rule - The emf around a closed curve  $C$  (or a curve that can be completed) is proportional to the change in flux through the surface  $S$  bounded by  $C$ .

$$\text{emf} = - \frac{d\Phi_m}{dt}$$

$\Rightarrow$  The flux rule is more powerful than the emf integral.

Ex Compute emf between ends of half circle moving with velocity  $v_0$  through constant field  $B_0$ .



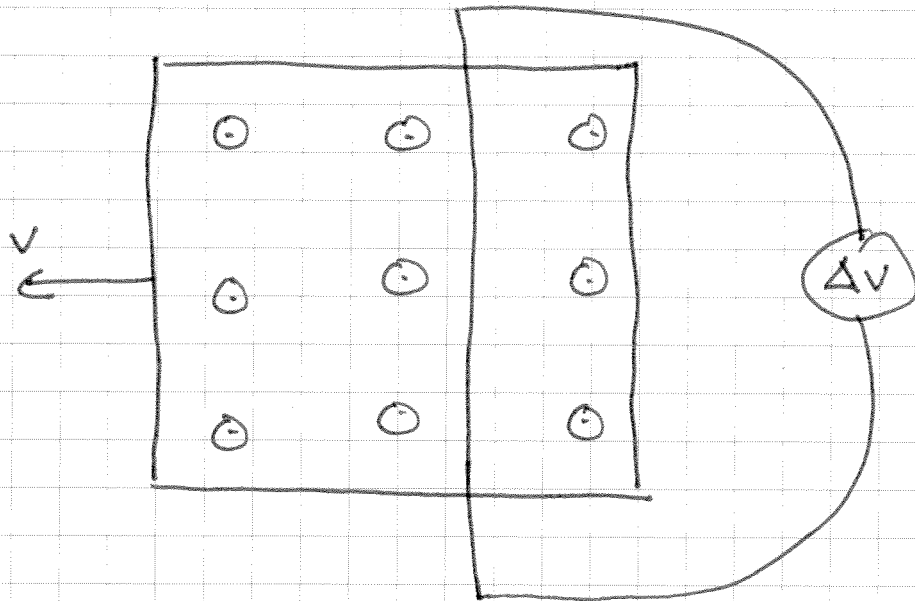
The flux through the surface above is

$$\Phi_m = B_0 2R v_0 (t - t_0)$$

The emf using the flux rule

$$\text{emf} = -\frac{d\Phi_m}{dt} = -2R v_0 B_0$$

Now consider the opposite experiment with our moving wire; move the magnet instead



$\Rightarrow$  Relativity says we should measure some  $\Delta V$ .

$$\text{emf} = \int_C (\vec{E} + \underbrace{\vec{v} \times \vec{B}}_0) \cdot d\vec{l}$$

$$\Rightarrow \vec{E} \neq 0$$

$\Rightarrow$  A changing magnetic field produces an electric field.

Faraday's Law The integral of the electric field around a closed path  $C$  is proportional to the time rate of change of the magnetic flux through a surface  $S$  bounded by  $C$ .

$$\oint_C \vec{E} \cdot d\vec{l} = - \frac{d\Phi_m}{dt} = - \frac{d}{dt} \int_S \vec{B} \cdot \hat{n} da$$

$\Rightarrow$  The path  $C$  is fixed.

$\Rightarrow$  The positive normal of  $S$  is given by the normal right hand convention

$\Rightarrow$  Special case of flux rule.

$\Rightarrow$  The direction of the electric field and therefore the direction of induced current, if  $C$  follows a conducting path follows direction from the convention for the positive normal.

Sorting out the normal directions gives Lenz' Law

Lenz' Law - The induced current flows in a direction s.t. the flux of the induced current through  $S$  opposes the changing flux through  $S$ .



## Differential Form

$$\oint_C \vec{E} \cdot d\vec{l} \stackrel{\text{Stokes}}{=} \int_S (\nabla \times \vec{E}) \cdot d\vec{l}$$
$$= - \frac{d}{dt} \int_S \vec{B} \cdot d\vec{a}$$

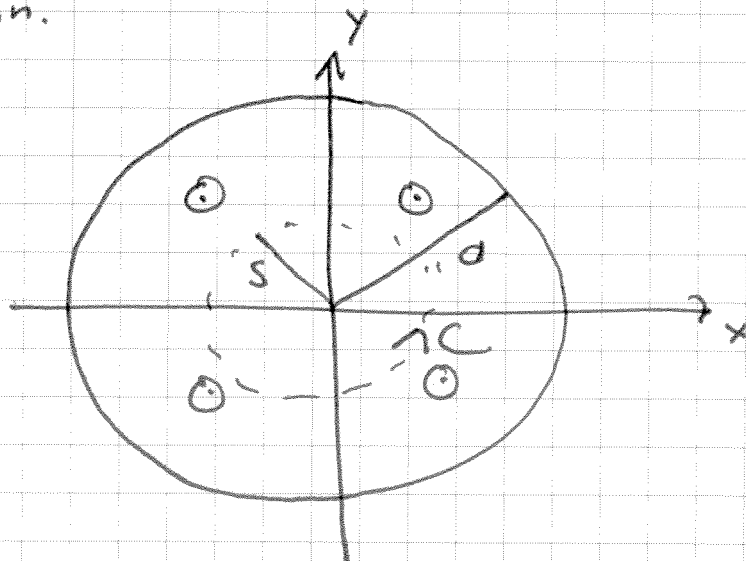
$$\nabla \times \vec{E} = - \frac{\partial \vec{B}}{\partial t} \quad \text{Faraday's Law}$$

⇒ Correct everywhere without restrictions

⇒ Does not require a loop of wire.

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Ex Consider a cylindrical region  $S \ll R$  containing a magnetic field  $\vec{B} = B_0 \frac{t}{\tau} \hat{z}$  where  $\tau$  is constant. Compute the electric field in the region.



B increasing in strength out of page

⇒ By RHR, positive normal to  $S$  out of page

⇒ The flux through a surface of radius  $s$  is

$$\begin{aligned}\Phi_m &= \oint \vec{B} \cdot d\vec{\sigma} = \pi s^2 B \\ &= \pi s^2 B_0 t / \tau\end{aligned}$$

Note sign because  $\hat{n} = \hat{z}$ .

### Faraday's Law

$$\oint_C \vec{E} \cdot d\vec{l} = - \frac{d\Phi_m}{dt} = - \frac{\pi s^2 B_0}{\tau}$$

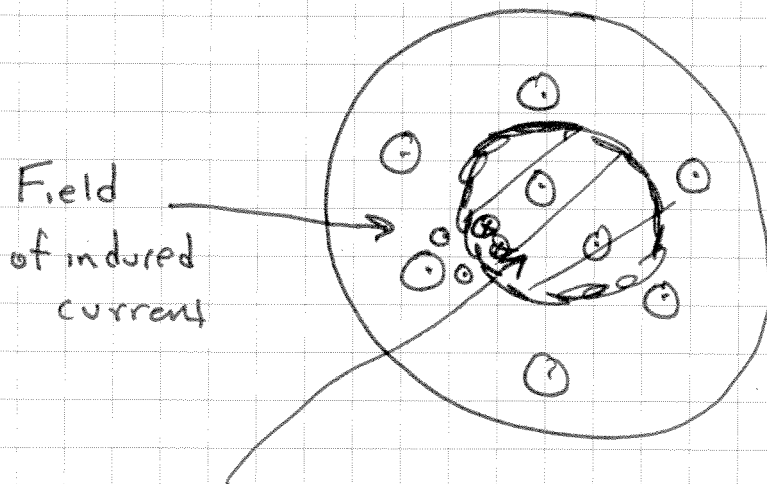
By symmetry, the electric field must be circular about the  $z$  axis,  $\vec{E} = E(s) \hat{\phi}$ .

Note ~~at~~ our chosen path has  $d\vec{l} = s d\phi \hat{\phi}$

$$\begin{aligned}\oint \vec{E} \cdot d\vec{l} &= \int_0^{2\pi} s d\phi E(s) = 2\pi s E \\ &= - \frac{d\Phi_m}{dt} = - \frac{\pi s^2 B_0}{\tau}\end{aligned}$$

$$\vec{E} = -\frac{s B_0}{2\pi} \hat{\phi}$$

⇒ Check direction using Lenz' Law. The electric field must be in the direction of the current that would flow in a conducting loop placed in the ~~electric~~ field.



Flux increasing out of the page, so induced current produces flux into the page, induced current must flow clockwise by RHR.

$$\text{Clockwise} = -\hat{\phi} \quad \checkmark$$