

Laplace's Egn (Cartesian)

Many times we will solve for the potential for systems with substantial charge free areas. In these regions the potential obeys Laplace's equation.

$$\nabla^2 V = 0$$

We will work out the solution to Laplace's eqn in cartesian, cylindrical, and spherical coordinates.

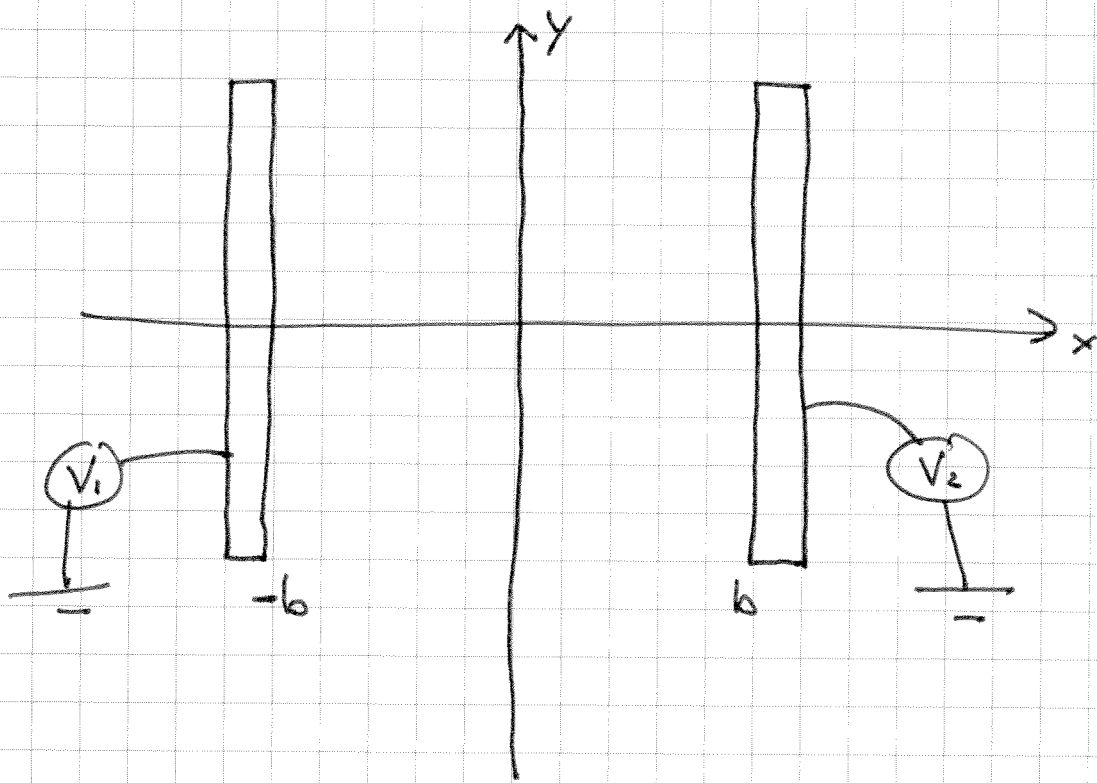
Cartesian Laplace's Egn

$$\nabla^2 V = \frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2} = 0$$

1D - The potential is constant in the y and z directions and only changes in x direction

$$V(\vec{r}) = V(x)$$

E_x Field between two conducting plates held at V_1 and V_2 .



Since $\rho = 0$ between the plates,

$$\nabla^2 V = \frac{\partial^2 V}{\partial x^2} = 0$$

Solutions $V=1$, $V=x$

General Solution

$$V(x) = a_1 + a_2 x$$

$$a_i = \text{constant}$$

Apply Boundary Conditions

$$V(b) = V_2 = a_1 + a_2 b$$

$$V(-b) = V_1 = a_1 - a_2 b$$

$$V_1 + V_2 = 2a_1$$

$$a_1 = \frac{V_1 + V_2}{2}$$

$$V_2 - V_1 = 2a_2 b$$

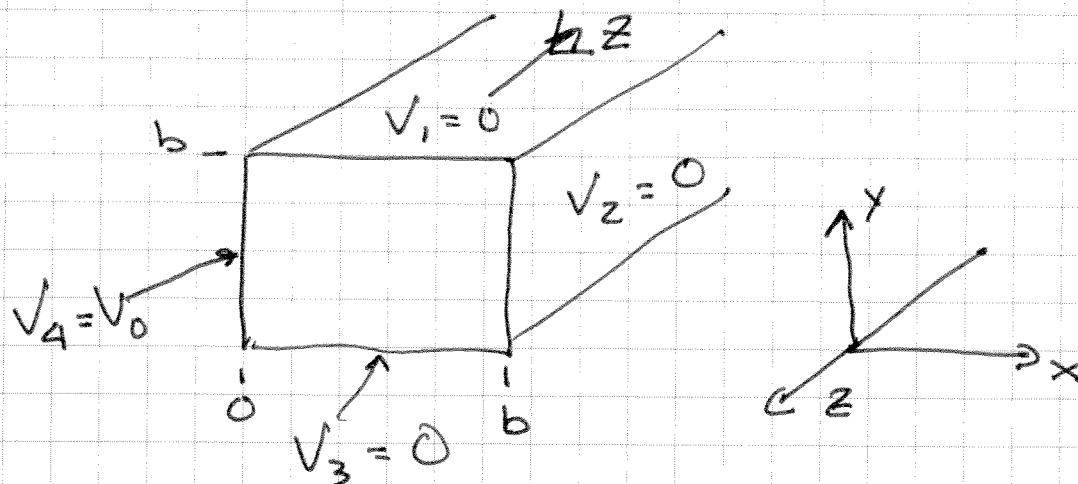
$$a_2 = \frac{V_2 - V_1}{2b}$$

Solution

$$V(x, y, z) = \frac{V_1 + V_2}{2} + \frac{V_2 - V_1}{2b} x$$

2D Cartesian

Potential does not change
in z direction \Rightarrow infinitely long channel.



Boundary Conditions (example)

$$V(b, y, z) = V_2 = 0$$

$$V(0, y, z) = V_4 = V_0$$

$$V(x, 0, z) = V_3 = 0$$

$$V(x, b, z) = V_1 = 0$$

Laplace's Egn

$$\nabla^2 V = \frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} = 0$$

With my choice

$$\frac{1}{X(x)} \frac{d^2 X(x)}{dx^2} = \kappa^2$$

$$\frac{1}{Y(y)} \frac{d^2 Y(y)}{dy^2} = -\kappa^2$$

$$\frac{d^2 X(x)}{dx^2} - \kappa^2 X(x) = 0$$

$$\frac{d^2 Y(y)}{dy^2} + \kappa^2 Y(y) = 0$$

Solutions

$$X(x) = \begin{matrix} e^{\kappa x} \\ e^{-\kappa x} \end{matrix}$$

$$Y(y) = \begin{matrix} \sin \kappa y \\ \cos \kappa y \end{matrix}$$

The boundary condition may not allow all values of κ .

Impose y Boundary Condition

$$Y(0) = 0 \quad Y(b) = 0$$

$$Y(0) = 0 \Rightarrow Y(y) = \sin \kappa x$$

$$Y(b) = 0 \Rightarrow \sin kb = 0$$

$$\Rightarrow kb = n\pi$$

$$\Rightarrow k = \frac{n\pi}{b} \equiv k_n$$

where n integer.

At this point any ~~$e^{\pm k_n x} \sin k_n y$~~

$$e^{\pm k_n x} \sin k_n y$$

satisfies Laplace's eqn and the y boundary condition.

The general solution is

$$V(x, y) = \sum_{n=1}^{\infty} \sin k_n y (c_n^+ e^{k_n x} + c_n^- e^{-k_n x})$$

c_n^{\pm} constants to be determined

Impose Boundary Conditions

$$V(0, y) = V_0 = \sum_{n=1}^{\infty} (c_n^+ + c_n^-) \sin k_n y$$

$$V(b, y) = 0 = \sum_{n=1}^{\infty} (c_n^+ e^{k_n b} + c_n^- e^{-k_n b}) \sin k_n y$$

To find c_n^{\pm} we need methods to pick these sums apart.

General Properties of Solution to Laplace's Eqn

Completeness - The set $\{f_n\}$ is complete if any function $f(x)$ satisfying $\nabla^2 f = 0$ and the boundary conditions can be written as a linear combination of $\{f_n\}$

$$f(x) = \sum_{n=1}^{\infty} a_n f_n(x)$$

where a_i are constants

We already implicitly used completeness to write the general solution as a sum.

Orthogonality

$$\int_0^b \sin k_n x \sin k_m x dx = \begin{cases} 0 & \text{if } m \neq n \\ \frac{b}{2} & \text{if } m = n \end{cases}$$
$$= \frac{b}{2} \delta_{mn}$$

Kronecker Delta

$$\delta_{mn} = \begin{cases} 0 & m \neq n \\ 1 & m = n \end{cases}$$

\Rightarrow The functions $\{\sin k_n x\}$ behave as orthogonal vectors.

\Rightarrow $\{\sin k_n x\}$ are linearly independent

\Rightarrow We cannot make one $\sin k_n x$ out of combinations of the other.

Fourier's Trick

To extract a_m from the series

$$f(x) = \sum_n a_n \sin k_n x = \sum_n a_n f_n(x)$$

Multiply by $f_m(x)$ and integrate,
 $f_m(x) = \sin k_m x$.

$$\int_0^b f_m(x) f(x) dx = \int_0^b f(x) \sin k_m x dx$$

$$= \sum_n \int_0^b a_n \sin k_n x \sin k_m x dx$$

$$= \sum_n a_n \frac{b}{2} \delta_{nm} = \frac{b}{2} a_m$$

$$a_m = \frac{2}{b} \int_0^b f(x) \sin k_m x dx$$

For our channel,

$$V(b, y) = 0 = \sum_{n=1}^{\infty} (c_n^+ e^{k_n b} + c_n^- e^{-k_n b}) \sin k_n y$$

Since $\sin k_n y$ are linearly independent, the coefficients of each must be zero.

$$c_n^+ e^{k_n b} + c_n^- e^{-k_n b} = 0$$

$$c_n^- = -c_n^+ e^{2k_n b}$$

$$V(0, y) = \sum_{n=1}^{\infty} c_n^+ (1 - e^{2k_n b}) \sin k_n y = V_0$$

Fourier's Trick Multiply by $\sin k_m y$ and integrate

$$\int_0^b V_0 \sin k_m y dy = \sum_{n=1}^{\infty} c_n^+ (1 - e^{2k_n b}) \times \underbrace{\int_0^b \sin k_m y \sin k_n y dy}_{\frac{b}{2} \delta_{nm}}$$

$$\int_0^b V_0 \sin k_m y \, dy = \sum_{n=1}^{\infty} C_n^+ (1 - e^{2k_n b}) \frac{b}{2} \sigma_{nm}$$

$$= C_m^+ \frac{b}{2} (1 - e^{2k_m b})$$

$$\int_0^b V_0 \sin k_m y \, dy = V_0 \int_0^b \sin \frac{m\pi y}{b} \, dy$$

$$= V_0 \frac{b}{m\pi} \int_0^{m\pi} \sin u \, du$$

$$= \frac{V_0 b}{m\pi} (-\cos u) \Big|_0^{m\pi}$$

$$= \begin{cases} 0 & m \text{ even} \\ \frac{2V_0 b}{m\pi} & m \text{ odd} \end{cases}$$

$$= C_m^+ \frac{b}{2} (1 - e^{2k_m b})$$

$$C_m^+ = \begin{cases} 0 & m \text{ even} \\ \frac{4V_0}{m\pi (1 - e^{2k_m b})} & m \text{ odd} \end{cases}$$

Full Solution

$$V(x,y) = \sum_n (c_n^+ e^{k_n x} + c_n^- e^{-k_n x}) \sin k_n y$$
$$= \sum_{n \text{ odd}} \frac{4V_0 \sin k_n y}{n\pi(1-e^{-2k_n b})} (e^{k_n x} - e^{2k_n b} e^{-k_n x})$$

* From this $\vec{E} = -\nabla V$

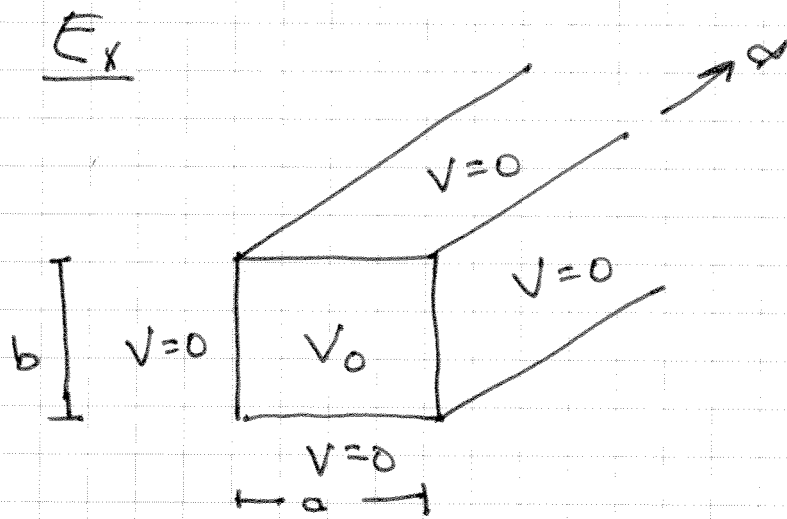
* The charge density on ~~one~~ of the conducting faces can be found as

$$\sigma = \epsilon_0 E_n = -\epsilon_0 \frac{\partial V}{\partial n}$$

E_x Charge on plate in $x-z$ plane

$$\sigma = -\epsilon_0 \left. \frac{\partial V}{\partial x} \right|_{x=0}$$

$$= \sum_{n \text{ odd}} \frac{-4\epsilon_0 V_0 \sin k_n y}{n\pi(1-e^{2k_n b})} (k_n + k_n e^{2k_n b})$$



Semi-infinite channel with one end held at V_0 and sides grounded.

Solutions (after boundary conditions)

$$X(x) = \sin k_n x \quad k_n = \frac{n\pi}{a}$$

$$Y(y) = \sin l_m y \quad l_m = \frac{m\pi}{b}$$

$$Z(z) = e^{wz}, e^{-wz}$$

\Rightarrow Discard e^{wz} because blows up at ∞

$$V(x, y, z) = \sum_{mn} A_{mn} \sin k_n x \sin l_m y e^{-(k_n^2 + l_m^2)^{1/2} z}$$