

Section - Systems of Particles

Consider a system of N (possibly interacting) particles with mass m_i , position \vec{r}_i , and velocity $\dot{\vec{r}}_i = \vec{v}_i$. Now we can describe anything.

Choice of system arbitrary.

Total Mass of System (m) $m = \sum_i m_i$

Center of Mass (\vec{r}_{cm}) $\vec{r}_{cm} = \frac{\sum m_i \vec{r}_i}{m}$

Total Momentum of System (\vec{p}) $\vec{p} = \sum \vec{p}_i = \sum m_i \vec{v}_i$

Velocity of Center of Mass (\vec{v}_{cm})

$$\vec{v}_{cm} = \dot{\vec{r}}_{cm} = \frac{1}{m} \sum m_i \dot{\vec{r}}_i$$

$$m \vec{v}_{cm} = \sum m_i \dot{\vec{r}}_i = \sum m_i \vec{v}_i = \vec{p}$$

The total momentum of the system is the same as a point mass, with total mass of the system moving at the velocity of the center of mass.

Total Force on Particle i

Net External Forces \vec{F}_i

Internal Forces $\sum_j \vec{F}_{ji}$

Newton II -

$$\dot{\vec{p}}_i = \vec{F}_i + \sum_j \vec{F}_{ji}$$

Sum Over all Particles

$$\sum_i \dot{\vec{p}}_i = \sum_i \vec{F}_i + \sum_i \sum_j \vec{F}_{ji}$$

Newton III - The internal force come in pairs $\vec{F}_{ji} = -\vec{F}_{ij}$, so the second sum is zero.

$$\sum_i \dot{\vec{p}}_i = \sum_i \vec{F}_i$$

Time Rate of Change of Total Momentum

$$\dot{\vec{p}} = m \dot{\vec{v}}_{cm} = m \vec{a}_{cm} = \sum_i \dot{\vec{p}}_i = \sum_i \vec{F}_i$$

The acceleration of the center of mass is the same as the acceleration of a point mass, with the total mass of the system, acted on by the sum of the forces acting on the system.

Body Forces Force of Gravity

$$\vec{F}_i = m_i \vec{g}$$

$$\sum \vec{F}_i = \sum m_i \vec{g} = m \vec{g} = m \vec{a}_{cm}$$

If no external forces,

$$\sum \vec{F}_i = 0 = \dot{\vec{P}} \quad \Rightarrow \quad \vec{P} = \text{constant}$$

Total Momentum is Conserved.

Section - Angular Momentum

Angular Momentum (\vec{L}) about the origin.

$$\vec{L} = \sum_i \vec{r}_i \times \vec{p}_i = \sum_i (\vec{r}_i \times m \vec{v}_i)$$

$$\frac{d\vec{L}}{dt} = \sum_i \left(\dot{\vec{r}}_i \times m \vec{v}_i + \vec{r}_i \times m \dot{\vec{v}}_i \right)$$

$$\vec{v}_i \times \vec{v}_i = 0$$

$$\dot{\vec{v}}_i = \vec{a}_i$$

*

$$m \vec{a}_i = \vec{F}_i + \sum_j \vec{F}_{ij}$$

$$\frac{d\vec{L}}{dt} = \sum_i \vec{r}_i \times \left(\vec{F}_i + \sum_j \vec{F}_{ij} \right)$$

$$= \sum_i \vec{r}_i \times \vec{F}_i + \sum_i \sum_j \vec{r}_i \times \vec{F}_{ij}$$

Newton III again - In the sum, we will

have pairs

$$\vec{r}_i \times \vec{F}_{ij} + \vec{r}_j \times \vec{F}_{ji}$$

$$\begin{aligned} & \vec{r}_i \times \vec{F}_{ij} + \vec{r}_j \times \vec{F}_{ji} \\ & \vec{F}_{ji} = -\vec{F}_{ij} \\ & (\vec{r}_j - \vec{r}_i) \times \vec{F}_{ij} \end{aligned}$$

The displacement vector from particle i to particle j is $\vec{r}_{ij} = \vec{r}_j - \vec{r}_i$

If \vec{F}_{ij} is central, $\vec{F}_{ij} \parallel \vec{r}_{ij}$ and the cross product is zero.

$$\frac{d\vec{L}}{dt} = \sum \vec{r}_i \times \vec{F}_i \equiv \vec{N} \quad \text{Torque}$$

If system is isolated $\vec{L} = 0$, $\vec{L} = \text{constant}$
angular momentum is conserved.

The conservation of momentum, angular momentum, and energy works in all known cases

SectionRotation about center of mass

Write particles position and velocity about the center of mass

$$\vec{r}'_i + \vec{r}_{cm} = \vec{r}_i$$

$$\vec{v}'_i + \vec{v}_{cm} = \vec{v}_i$$

$$\vec{L} = \sum_i \vec{r}_i \times \vec{p}_i = \sum_i (\vec{r}'_i + \vec{r}_{cm}) \times m_i (\vec{v}'_i + \vec{v}_{cm})$$

$$= \left(\sum_i m_i \right) \vec{r}_{cm} \times \vec{v}_{cm} + \vec{r}_{cm} \times \sum_i m_i \vec{v}'_i$$

$$+ \left(\sum_i m_i \vec{r}'_i \right) \times \vec{v}_{cm} + \sum_i \vec{r}'_i \times m_i \vec{v}'_i$$

$$\sum m_i \vec{r}'_i + m_i \vec{r}_{cm} = \sum m_i \vec{r}_i$$

$$\underbrace{\sum m_i \vec{r}'_i}_0 + m \vec{r}_{cm} = m \vec{r}_{cm}$$

Likewise $\sum m_i \vec{v}'_i = 0$

$$\vec{L} = \underbrace{m \vec{r}_{cm} \times \vec{v}_{cm}}_{\text{Angular Momentum of center of mass about origin.}} + \underbrace{\sum \vec{r}'_i \times m_i \vec{v}'_i}_{\text{Angular Momentum of system about center of mass.}}$$

Kinetic Energy (Likewise)

$$T = \sum_i \frac{1}{2} m_i \vec{v}_i \cdot \vec{v}_i$$

$$= \frac{1}{2} m v_{cm}^2 + \sum_i \frac{1}{2} m_i v_i'^2$$

Section - Reduced Mass

Consider an isolated two particle system with CM at origin.

$$m_1 \vec{r}_1' + m_2 \vec{r}_2' = 0 \quad \leftarrow \text{displacement from CM}$$

$$\vec{r}_2' = -\frac{m_1}{m_2} \vec{r}_1'$$

EOM of particle 1

$$m_1 \frac{d^2 \vec{r}_1'}{dt^2} = \vec{F}_{21} = f(r_{21}) \hat{r}_{21}$$

~~Suppose \vec{F}_1 is on i~~

Displacement Vector

$$\begin{aligned} \vec{r}_{21} &= \vec{r}_1' - \vec{r}_2' \\ &= \vec{r}_1' \left(1 + \frac{m_1}{m_2} \right) \end{aligned}$$

$$\frac{\vec{r}_{21}}{1 + \frac{m_1}{m_2}} = \vec{r}_1'$$

$$\frac{m_1}{1 + \frac{m_1}{m_2}} \frac{d^2 \vec{r}_{21}}{dt^2} = f(r_{21}) \hat{r}_{21}$$

Define reduced mass

$$\mu = \frac{m_1 m_2}{m_1 + m_2}$$

$$\mu \frac{d^2 \vec{r}_{z_1}}{dt^2} = f(r_{z_1}) \hat{r}_{z_1} \quad \leftarrow \text{orbits } z$$

The effect of having two particles instead of a fixed center of force is to replace the inertial mass m_1 with μ .

If f is gravity, $f = -\frac{m_1 m_2 G}{r^2}$

not $\frac{\mu m_2 G}{r^2}$

So we can go back and fix all our orbit equations. For example

$$\tau^2 = \frac{4\pi^2}{GM_0} a^3 \Rightarrow \tau^2 = \frac{4\pi^2}{G(M_0 + m)} a^3$$

Section 1
Lecture 3/29 - Collisions

In a collision between N particles, the forces are very short range and generally unknown. We observe the masses m_i and the momenta \vec{p}_i before and after collision. In nuclear collision, the m_i may change.

Momentum Conservation

$$\underbrace{\sum \vec{p}_i^I}_{\text{initial momentums}} = \underbrace{\sum \vec{p}_i^F}_{\text{final momentums}}$$

Energy Conservation

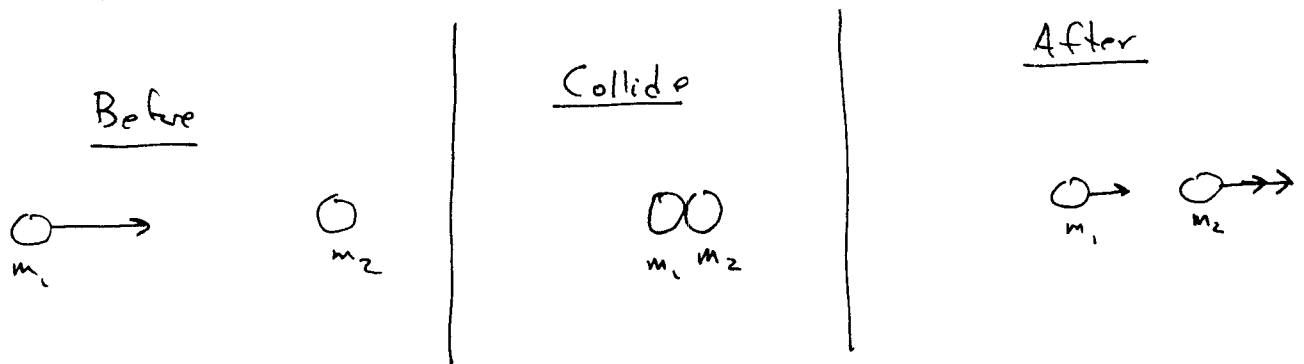
$$\sum_i \frac{(\vec{p}_i^I)^2}{2m_i} = \sum_i \frac{(\vec{p}_i^F)^2}{2m_i} + Q$$

$Q \equiv$ Energy lost to environment in collision. If $Q > 0$ the particles have less energy after the collision.

Section 2 - Direct Collisions

Back to OP III.

Suppose the particles ^{before and} after the collision travel on a single line. \Rightarrow Direct Collision (Head on collision)



Since all in a line, we don't have to use vectors anymore. They all have the same unit vector.

~~To further simplify, work in a coordinate system where m_2 is at rest, and only two particles collide.~~

Momentum Conservation (p have signs)

$$P_2^I + P_1^I = P_1^F + P_2^F$$

Energy Conservation

$$\frac{(P_2^I)^2}{2m_2} + \frac{(P_1^I)^2}{2m_1} = \frac{(P_1^F)^2}{2m_1} + \frac{(P_2^F)^2}{2m_2} + Q$$

Relative Velocity

$$V^I = |v_2^I - v_1^I| \quad \text{Before Collision}$$

$$V^F = |v_2^F - v_1^F| \quad \text{After Collision}$$

Coefficient of Restitution

(How much velocity is restored after collision?)

$$\epsilon = \frac{V^F}{V^I}$$

Elastic Collision - No energy lost or gained from environment, $Q=0$ and $\epsilon=1$.

PF

$$m_1 v_1^I + m_2 v_2^I = m_1 v_1^F + m_2 v_2^F$$

If elastic, $\epsilon=1$

$$v_2^F - v_1^F = v_1^I - v_2^I$$
$$(m_1 v_1^I + m_2 v_2^I)(v_1^I - v_2^I) = (m_1 v_1^F + m_2 v_2^F)(v_2^F - v_1^F)$$
$$m_1 (v_1^I)^2 + m_2 (v_2^I)^2 + (m_1 + m_2)(v_1^I v_2^I) = -m_1 v_1^F v_2^F$$

General Final Velocities

$$v_1^F = \frac{(m_1 - \epsilon m_2) v_1^I + (m_2 + \epsilon m_1) v_2^I}{m_1 + m_2}$$

$$v_2^F = \frac{(m_1 + \epsilon m_2) v_1^I + (m_2 - \epsilon m_1) v_2^I}{m_1 + m_2}$$

Totally Inelastic $\epsilon = 0$

$$v_1^F = \frac{m_1 v_1^I + m_2 v_2^I}{m_1 + m_2} = v_{cm}$$

$$= v_2^F$$

The particles move with the same velocity.

Totally Elastic $\epsilon = 1$ AND $m_1 = m_2$

$$v_1^F = v_2$$

$$v_2^F = v_1$$

Exchange velocities.

Energy Loss

$$Q = \frac{1}{2} \mu (v_I)^2 (1 - \epsilon^2)$$

Left as homework problem.

\Rightarrow How large can Q be?

$$Q_{\max} = \frac{1}{2} \mu (v_I)^2$$

Recall The kinetic energy with respect to center of mass can be separated as

$$T = \frac{1}{2} M v_{cm}^2 + \sum \frac{1}{2} m_i v_i'^2$$

\leftarrow velocity with respect to cm.

With about a page of algebra, you can show

$$Q_{\max} = \sum \frac{1}{2} m_i v_i'^2$$

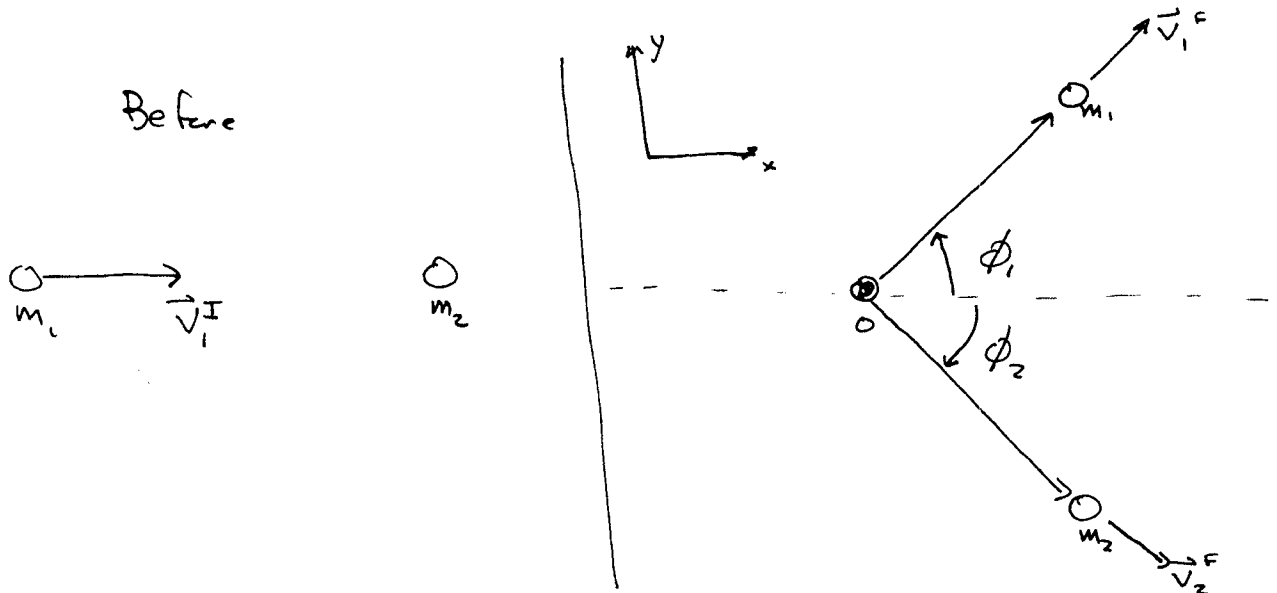
The energy of the motion of the center of mass cannot be converted to another form. Required to conserve momentum,

\Rightarrow Q does not depend on coordinate system, since depends relative velocity.

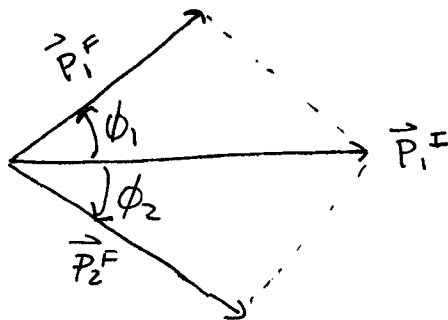
Section

Collisions in Two Dimensions

Select a coordinate system, the lab system,
where m_2 is at rest before collision.



Momentum Conservation



x-component

$$p_1^I = p_1^F \cos \phi_1 + p_2^F \cos \phi_2$$

y-component

$$0 = p_1^F \sin \phi_1 - p_2^F \sin \phi_2$$

Energy

$$\frac{(P_1^I)^2}{2m_1} = \frac{(P_1^F)^2}{2m_1} + \frac{(P_2^F)^2}{2m_2} + Q$$

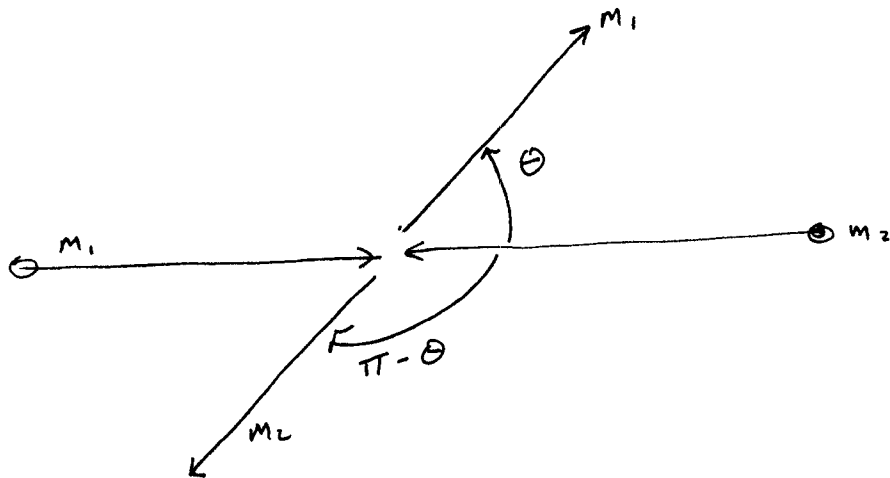
How does all the coefficient of restitution jazz work here?

Consider the collision in CM coordinate system, center of mass fixed at origin.

$$\vec{v}_1 = \vec{v}_{cm} + \vec{v}_1' \quad \vec{v}_2 = \vec{v}_{cm} + \vec{v}_2'$$

Total Momentum $\vec{P} = M\vec{v}_{cm} \Rightarrow \vec{P} = 0$

$$\vec{P}_1^{I'} + \vec{P}_2^{I'} = 0 = \vec{P}_1^{F'} + \vec{P}_2^{F'}$$



Particles move in line after collision. Can define relative velocity v^I and v^F , coefficient of restitution. It all still works in CM, ONLY. Note whether particles exchange velocity depends on scattering angle.

Relation between Lab and CM scattering angles.



x component

$$v_{iF} \cos \phi_1 = v_{cm} + v_{iF}' \cos \theta$$

y component

$$v_{iF} \sin \phi_1 = v_{iF}' \sin \theta$$

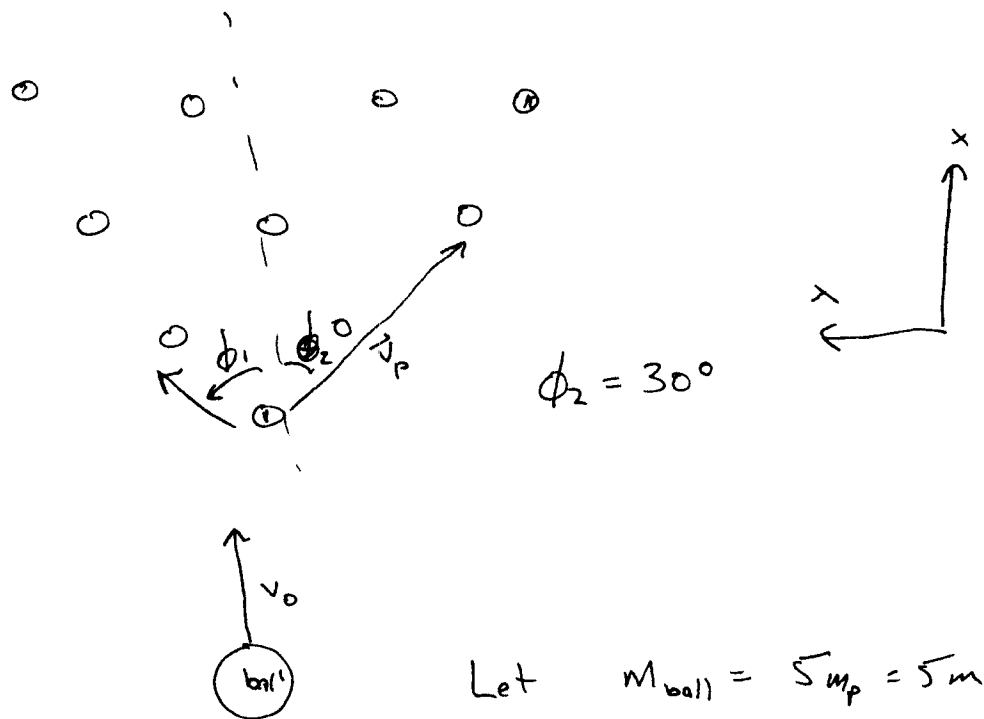
$$\tan \phi_1 = \frac{v_{iF}' \sin \theta}{v_{cm} + v_{iF}' \cos \theta} = \frac{\sin \theta}{\gamma + \cos \theta}$$

$$\gamma = v_{cm} / v_{iF}'$$

For an elastic collision, $\gamma = m_1 / m_2$

Through a bit of algebra.

Section - Example 1/10 Split



$$\vec{P}_{\text{ball}}^I = \vec{P}_{\text{ball}}^F + \vec{P}_{\text{pin}}^F$$

$$\frac{1}{2} M_{\text{ball}} v_{\text{ball}}^I{}^2 = \frac{1}{2} M_{\text{ball}} v_{\text{ball}}^F{}^2 + \frac{1}{2} M_{\text{pin}} v_{\text{pin}}^F{}^2 + Q$$

Let $Q=0$, approximate collision as elastic, no rotation.

x-component momentum

$$5m v_0 = 5m v_b^F \cos \phi_1 + m v_p^F \cos 30^\circ \quad (1)$$

y-component momentum

$$0 = 5m v_b^F \sin \phi_1 - m v_p^F \sin 30^\circ \quad (2)$$

Energy

$$5V_0^2 = 5V_b^{F2} + V_p^{F2}$$

$$V_0^2 - \frac{1}{5} V_p^{F2} = V_b^{F2} \quad (3)$$

Manipulate

$$V_0 - \frac{1}{5} V_p^F \cos 30^\circ = V_b^F \cos \phi_1 \quad (4)$$

$$\frac{1}{5} V_p^F \sin 30^\circ = V_b^F \sin \phi_1 \quad (5)$$

$$\underline{(4)^2 + (5)^2 = (6)}$$

$$V_0^2 - \frac{2}{5} V_p^F V_b^F \cos 30^\circ + \frac{1}{25} V_p^{F2} \cos^2 30^\circ + \frac{1}{25} V_p^{F2} \sin^2 30^\circ = V_b^{F2} \quad (6)$$

(6) - Energy

$$-\frac{2}{5} V_p^F V_b^F \cos 30^\circ + \frac{1}{25} V_p^{F2} + \frac{1}{5} V_p^{F2} = 0$$

$$-10 V_0 \cos 30^\circ + 6 V_p^F = 0$$

$$V_p^F = \frac{10}{6} V_0 \cos 30^\circ = 1.44 V_0$$

Angle (5)/(4)

$$\begin{aligned}\tan \phi_1 &= \frac{\frac{1}{5} V_P^F \sin 30^\circ}{V_0 - \frac{1}{5} V_P^F \cos 30^\circ} \\ &= \frac{1}{\frac{5V_0}{V_P^F \sin 30^\circ} - \frac{\cos 30^\circ}{\sin 30^\circ}} \\ &= 0.19\end{aligned}$$

$$\phi_1 \approx 10^\circ$$

Velocity of Ball

$$V_b^F = \sqrt{V_0^2 - \frac{1}{5} V_P^F^2} = 0.77V_0$$

$$\vec{V}_b^F = 0.77V_0 (\cos 10, \sin 10, 0)$$

$$\vec{V}_P^F = 1.44V_0 (\cos 30, -\sin 30, 0)$$