

Example Test 2

Justin's Version

CO₂ Molecule

A CO₂ molecule can be modeled as three masses connected by 2 springs of spring constant k . Let the mass of the carbon atom be $m_c = 3m$ and the mass of the oxygen atoms each be $m_o = 4m$ (even though these numbers may not look right, the ratio is correct). Allow only motions along the molecular axis (back and forth, not up and down).

- a) Find the Lagrangian of this system.
- b) Find the equations of motion of the system.
- c) Find the Hamiltonian of the system.
- d) Find the frequencies of the system.
- e) Find the normal modes of the system. Are all these vibrations?

■ Solution

a)

$$L = T - V$$

$$= \frac{1}{2} (m_O \dot{x}_1^2 + m_c \dot{x}_2^2 + m_O \dot{x}_3^2) - \frac{k}{2} ((x_2 - x_3)^2 + (x_3 - x_2)^2)$$

b)

Use $\frac{\partial L}{\partial q} - \frac{d}{dt} \frac{\partial L}{\partial \dot{q}}$ to find the following.

$$m_O \ddot{x}_1 + k x_1 - k x_2 = 0$$

$$-k x_1 + m_c \ddot{x}_2 + 2k x_2 - k x_3 = 0$$

$$-k x_2 + m_O \ddot{x}_3 + k x_3 = 0$$

c)

$$H = T + V$$

$$= \frac{1}{2} (m_O \dot{x}_1^2 + m_c \dot{x}_2^2 + m_O \dot{x}_3^2) + \frac{k}{2} ((x_2 - x_3)^2 + (x_3 - x_2)^2)$$

Put this into the form $H(q_i, p_i)$.

$$H = \frac{1}{2} \left(\frac{p_1^2}{m_O} + \frac{p_2^2}{m_c} + \frac{p_3^2}{m_O} \right) + \frac{k}{2} ((x_2 - x_3)^2 + (x_3 - x_2)^2)$$

d)

Place the equations of motion in a matrix. Also assume that the motions are periodic.

$$\begin{pmatrix} k - m_O \omega^2 & -k & 0 \\ -k & 2k - m_c \omega^2 & -k \\ 0 & -k & k - m_O \omega^2 \end{pmatrix} \begin{pmatrix} a_1 \\ a_2 \\ a_3 \end{pmatrix} = 0$$

Giving the following secular equation.

$$\text{Det} \left[\begin{pmatrix} k - m_O \omega^2 & -k & 0 \\ -k & 2k - m_c \omega^2 & -k \\ 0 & -k & k - m_O \omega^2 \end{pmatrix} \right]$$

$$-k^2 \omega^2 m_c - 2k^2 \omega^2 m_O + 2k \omega^4 m_c m_O + 2k \omega^4 m_O^2 - \omega^6 m_c m_O^2$$

Solve the secular equation for ω^2 . This gives 3 eigenvalues. If you get confused how to solve a cubic equation look to your mathematical handbook. The other trick is to notice that all terms have at least ω^2 in them. That means $\omega = 0$ is clearly an eigenvalue. You can divide by ω^2 and have a quadratic you can solve.

$$\omega_1 = 0$$

$$\omega_2 = \sqrt{\frac{k}{m_o}}$$

$$\omega_3 = \sqrt{\frac{k}{m_o} + 2 \frac{k}{m_c}}$$

e)

Any mode with zero frequency is suspect. Plug into our earlier matrix to see that $a_1, a_2, a_3 = 0$.

Try ω_2

$$\begin{pmatrix} k - m_o \frac{k}{m_o} & -k & 0 \\ -k & 2k - m_c \frac{k}{m_o} & -k \\ 0 & -k & k - m_o \frac{k}{m_o} \end{pmatrix} \cdot \begin{pmatrix} a_1 \\ a_2 \\ a_3 \end{pmatrix}$$

$$\text{MatrixForm} \left[\begin{pmatrix} k - m_o \frac{k}{m_o} & -k & 0 \\ -k & 2k - m_c \frac{k}{m_o} & -k \\ 0 & -k & k - m_o \frac{k}{m_o} \end{pmatrix} \cdot \begin{pmatrix} a_1 \\ a_2 \\ a_3 \end{pmatrix} \right]$$

$$\begin{pmatrix} -k a_2 \\ -k a_1 - k a_3 + a_2 \left(2k - \frac{k m_c}{m_o} \right) \\ -k a_2 \end{pmatrix}$$

For all elements of this array to equal zero $a_2 = 0$. From this you quickly see that $a_1 = -a_3$. Spectroscopists will call this an antisymmetric mode.

Now look for ω_3 .

$$\text{MatrixForm} \left[\begin{pmatrix} k - m_o \left(\frac{k}{m_o} + 2 \frac{k}{m_c} \right) & -k & 0 \\ -k & 2k - m_c \left(\frac{k}{m_o} + 2 \frac{k}{m_c} \right) & -k \\ 0 & -k & k - m_o \left(\frac{k}{m_o} + 2 \frac{k}{m_c} \right) \end{pmatrix} \cdot \begin{pmatrix} a_1 \\ a_2 \\ a_3 \end{pmatrix} \right]$$

$$\begin{pmatrix} -k a_2 + a_1 \left(k - m_o \left[\frac{2k}{m_c} + \frac{k}{m_o} \right] \right) \\ -k a_1 - k a_3 + a_2 \left(2k - m_c \left[\frac{2k}{m_c} + \frac{k}{m_o} \right] \right) \\ -k a_2 + a_3 \left(k - \left(\frac{2k}{m_c} + \frac{k}{m_o} \right) m_o \right) \end{pmatrix}$$

From the first and third element we see that $a_1 = a_3$, else they could not be equal zero. More algebra will show that $a_2 = -a_1(m_o/m_c)$.

Stability

A particle of mass m moves in the following potential. Is there a stable orbit? If so, what is the frequency, ω , of the small oscillatory motion?

1) $V = k x^{-2} e^x + c$

2) $V = A \cos(kx + \pi/4)$

■ Solution

1)

Take the first derivative to see if there is a min or max.

$$\frac{k_2 e^x}{x^2} - \frac{2 k_2 e^x}{x^3}$$

A bit of algebra later and we see that the first derivative has a zero at $x = 2$. Look to the second derivative to see if it is a max or min.

$$\frac{6 k_2 e^x}{x^4} - \frac{4 k_2 e^x}{x^3} + \frac{k_2 e^x}{x^2}$$

At $x = 2$ this is positive, showing a minimum. We can find the effective spring constant by looking at the value of the second derivative.

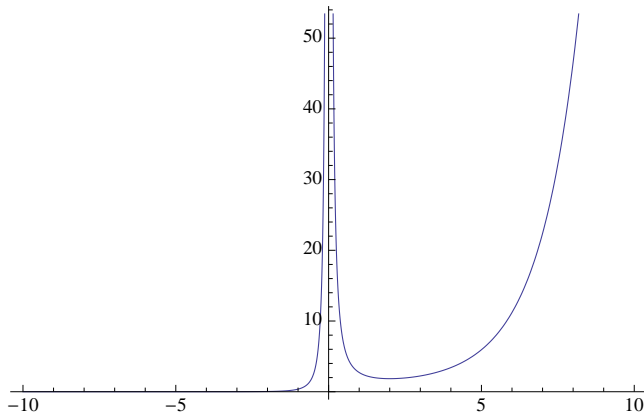
We do this and see the following.

$$k_{\text{eff}} = V'' = \frac{e^2 k}{8}$$

Frequency follows from this.

$$\omega = \sqrt{\frac{k_{\text{eff}}}{m}}$$

Plot $[x^{-2} e^x, \{x, -10, 10\}]$



2)

Zeros for this are easy to find. Were the function simply $\cos(x)$ the zeros would be at $x=0, \pi/k, 2\pi/k, n\pi/k$. For this function we shift by $\pi/4$, giving minima at $x = -\pi/4k, 3\pi/4k, (4n-1)\pi/4k$. Only half of these are minima ($3\pi/4k$). The others are maxima.

The frequency can be found the same way as before.

$$k_{\text{eff}} = A k^2$$

$$\omega = \sqrt{\frac{k_{\text{eff}}}{m}}$$

Double Atwood's Machine.

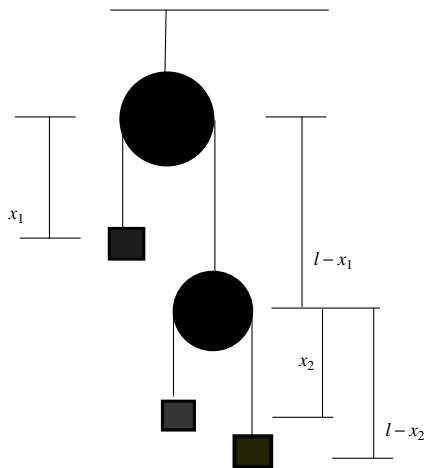
The following diagram shows a double Atwood's machine. Let each pulley be of mass M and radius r , each block be of mass m_i , and each string be of length l .

1)

Use the Lagrangian to find the equations of motion for this system.

2)

Finding a Hamiltonian in terms of q_i and p_i is difficult for the double Atwood's Machine. Why might that be?



■ Solution

1)

$$T = \frac{1}{2} \left(m_1 \dot{x}_1^2 + m_2 (\dot{x}_2 - \dot{x}_1)^2 + m_3 (-\dot{x}_2 - \dot{x}_1)^2 + M (-\dot{x}_1)^2 + I \frac{\dot{x}_1^2}{a^2} + I \frac{\dot{x}_2^2}{a^2} \right)$$

$$V = -m_1 g x_1 - M g(l - x_1) - m_2 g(l - x_1 + x_2) - m_3 g(l - x_1 + l - x_2)$$

$$L = \frac{1}{2} \left(m_1 \dot{x}_1^2 + m_2 (\dot{x}_2 - \dot{x}_1)^2 + m_3 (-\dot{x}_2 - \dot{x}_1)^2 + M (-\dot{x}_1)^2 + I \frac{\dot{x}_1^2}{a^2} + I \frac{\dot{x}_2^2}{a^2} \right) + m_1 g x_1 + M g(l - x_1) + m_2 g(l - x_1 + x_2) + m_3 g(l - x_1 + l - x_2)$$

Use $\frac{\partial L}{\partial q} = \frac{d}{dt} \frac{\partial L}{\partial \dot{q}}$ to find the equations of motion.

 x_1 :

$$(m_1 - M - m_2 - m_3) g = m_1 \ddot{x}_1 + m_2 (\ddot{x}_1 - \ddot{x}_2) + m_3 (\ddot{x}_2 + \ddot{x}_1) + M \ddot{x}_1 + I \frac{\ddot{x}_1}{a^2}$$

 x_2 :

$$m_2 g - m_3 g = m_2 (\ddot{x}_2 - \ddot{x}_1) + m_3 (\ddot{x}_2 + \ddot{x}_1) + I \frac{\ddot{x}_2}{a^2}$$

2)

While T is easy to write in terms of velocities, for this case it is harder to write in terms of momenta. $\frac{\partial L}{\partial \dot{x}_1} = p_1$ is not simply $m_1 \dot{x}_1$ as in other problems we have worked.

Forces of Constraint

A small, wet bar of soap of mass m can move about the inside of a hemispherical bowl of radius R .

- 1) Write the Lagrangian for this system.
- 2) What is the constraint for this system.
- 3) Write the equations of motion.
- 4) What is the normal force on the soap?

■ **Solution**

1)

$$L = \frac{m}{2} (\dot{x}^2 + \dot{y}^2 + \dot{z}^2) - m g z$$

$$= \frac{m}{2} (\dot{r}^2 + r^2 \dot{\theta}^2 + r^2 \dot{\phi}^2 \sin^2(\theta)) - m g r(1 - \cos(\theta))$$

ϕ must be constant, so it may be ignored.

2)

The constraint it for the soap to only move along the hemisphere.

$$r = R; \dot{r} = 0$$

$$f = r - R$$

3)

$$L = \frac{m}{2} (\dot{r}^2 + r^2 \dot{\theta}^2) - m g r(1 - \cos(\theta))$$

r :

$$m g \cos(\theta) + m r \dot{\theta}^2 + m \ddot{r} + \lambda_r = 0$$

we know that $\ddot{r} = 0$. So we can simplify

$$m g \cos(\theta) + m r \dot{\theta}^2 = \lambda_r$$

θ :

$$-m g \sin(\theta) + m r^2 \ddot{\theta} = 0$$

or

$$\ddot{\theta} = \frac{g}{r^2} \sin(\theta)$$

4)

The normal force is already done in part 3.

$$F_N = \lambda_r \frac{\partial f}{\partial r} = Q_r = m g \cos(\theta) + m r \dot{\theta}^2$$