

## Angular Momentum Matrices

The generalized angular momentum matrices

$J^2, J_z$  had quantum numbers

$$J = 0, \frac{1}{2}, 1, \frac{3}{2}, 2, \dots$$

The integer  $J$ 's have position space eigenfunctions and can be used to represent orbital angular momentum. The  $\frac{1}{2}$  integer  $J$ 's do not have eigenfunctions, but do represent a physical property, the intrinsic angular momentum or spin. We must use matrices to investigate  $J = \frac{1}{2}, \frac{3}{2}, \dots$

## Build Angular Momentum Matrices

$$\underline{J=0} \quad \Rightarrow \quad m=0$$

Basis States  $\{ |00\rangle \}$

(2)

$J^2$  and  $J_z$  are  $1 \times 1$  matrices

$$J^2 = (\langle 00 | J^2 | 00 \rangle) = (0)$$

$$J_z = (\langle 00 | J_z | 00 \rangle) = (0)$$

$$\underline{J = 1/2, \quad m = \pm 1/2}$$

Basis States  $\{ | \frac{1}{2} \frac{1}{2} \rangle, | \frac{1}{2} -\frac{1}{2} \rangle \}$

$$| + \rangle \equiv | \frac{1}{2} \frac{1}{2} \rangle$$

$$| - \rangle \equiv | \frac{1}{2} -\frac{1}{2} \rangle$$

$$J^2 | + \rangle = \frac{1}{2} (\frac{1}{2} + 1) \hbar^2 | + \rangle = J^2 | - \rangle$$

$$= \frac{3}{4} \hbar^2 | + \rangle$$

$$J_z | + \rangle = \frac{1}{2} \hbar | + \rangle$$

$$J_z | - \rangle = -\frac{1}{2} \hbar | - \rangle$$

3

A vector in the  $\{|+\rangle, |-\rangle\}$  basis is called a "spinor".

We will use the operators  $\vec{S}, S^2, S_z$  for spin  $1/2$ .

$$S^2 = J^2 \text{ etc.}$$

Find  $S^2, S_x, S_y, S_z$

$$S^2 = \begin{pmatrix} \langle + | S^2 | + \rangle & \langle + | S^2 | - \rangle \\ \langle - | S^2 | + \rangle & \langle - | S^2 | - \rangle \end{pmatrix}$$

$$= \frac{3}{4} \hbar^2 \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$S_z = \begin{pmatrix} \langle + | S_z | + \rangle & \langle + | S_z | - \rangle \\ \langle - | S_z | + \rangle & \langle - | S_z | - \rangle \end{pmatrix}$$

$$= \frac{1}{2} \hbar \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

To calculate  $S_x, S_y$  work on raising/lowering ops

$$\hat{S}_+ = \hat{S}_x + i \hat{S}_y \quad \hat{S}_- = \hat{S}_x - i \hat{S}_y$$

$$S_x = \frac{1}{2} (\hat{S}_+ + \hat{S}_-) \quad S_y = \frac{1}{2i} (\hat{S}_+ - \hat{S}_-)$$

Normalized Operators

$$\hbar \sqrt{j(j+1) - m(m+1)} |j, m+1\rangle = J_+ |j, m\rangle$$

$$\hbar \sqrt{j(j+1) - m(m-1)} |j, m-1\rangle = J_- |j, m\rangle$$

S<sub>0</sub>

$$S_+ |+\rangle = S_+ |\frac{1}{2} \frac{1}{2}\rangle$$

$$= 0$$

$$S_+ |-\rangle = S_+ |\frac{1}{2} -\frac{1}{2}\rangle$$

$$= \hbar \sqrt{\frac{1}{2}(\frac{1}{2}+1) - (-\frac{1}{2})(-\frac{1}{2}+1)} |\frac{1}{2} \frac{1}{2}\rangle$$

$\frac{3}{4} \quad + \frac{1}{4}$

(5)

$$S_+ |-\rangle = \hbar |+\rangle$$

$$\langle + | S_+ |-\rangle = \hbar \langle + | + \rangle = \hbar$$

$$\langle - | S_- |+\rangle = \hbar$$

$$S_x = \begin{pmatrix} \langle + | \frac{1}{2}(S_+ + S_-) | + \rangle & \langle + | \frac{1}{2}(S_+ + S_-) | - \rangle \\ \langle - | \frac{1}{2}(S_+ + S_-) | + \rangle & \langle - | \frac{1}{2}(S_+ + S_-) | - \rangle \end{pmatrix}$$

$$= \begin{pmatrix} 0 & \frac{1}{2} \langle + | S_+ | - \rangle \\ \frac{1}{2} \langle - | S_- | + \rangle & 0 \end{pmatrix}$$

$$= \frac{1}{2} \hbar \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

Likewise

$$S_y = \frac{1}{2} \hbar \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

6

## Pauli Spin Matrices

$$\hat{\sigma}_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$\hat{\sigma}_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

$$\hat{\sigma}_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$\hat{S}_x = \frac{1}{2}\hbar \hat{\sigma}_x$$

$$\hat{S}_y = \frac{1}{2}\hbar \hat{\sigma}_y$$

$$\hat{S}_z = \frac{1}{2}\hbar \hat{\sigma}_z$$

7

Let's try  $L_x, L_y, L_z$  for  $\ell=1$ .

$$L^2 = \ell(\ell+1)\hbar^2 = 2\hbar^2$$

$$L^2 = 2\hbar^2 \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

State space  $|1\rangle \equiv |11\rangle$   $|0\rangle \equiv |10\rangle$   $|-1\rangle \equiv |1-1\rangle$

$$L_z = \hbar \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix}$$

As before  $\hat{L}_x = \frac{1}{2}(\hat{L}_+ + \hat{L}_-)$

$$L_x = \begin{pmatrix} \langle 1|L_x|1\rangle & \langle 1|L_x|0\rangle & \langle 1|L_x|-1\rangle \\ \langle 0|L_x|1\rangle & \langle 0|L_x|0\rangle & \langle 0|L_x|-1\rangle \\ \langle -1|L_x|1\rangle & \langle -1|L_x|0\rangle & \langle -1|L_x|-1\rangle \end{pmatrix}$$

⑧

Only non-zero terms are elements where bra and ket 1 apart.

$$L_+ |0\rangle = L_+ |10\rangle = \hbar \sqrt{1(1+1) - 0(0+1)} |11\rangle \\ = \sqrt{2} \hbar |11\rangle$$

$$L_+ |-1\rangle = \hbar \sqrt{1(1+1) + 1(-1+1)} |0\rangle \\ = \sqrt{2} \hbar |0\rangle$$

$$\langle 1 | S_x | 0 \rangle = \frac{1}{2} \langle 1 | S_+ + S_- | 0 \rangle = \frac{\hbar}{\sqrt{2}}$$

$$\langle 0 | S_x | -1 \rangle = \frac{1}{2} \langle 0 | S_+ + S_- | -1 \rangle = \frac{\hbar}{\sqrt{2}}$$

$$L_x = \frac{\hbar}{\sqrt{2}} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}$$

Other elements can be found by Hermiticity