

# PHYS 4073 - Quantum Mechanics- Homework Set 6

Reading Assignment: No additional material.

Due at 5:45pm Monday October 25th in my box or at my office.

## Griffiths' Problems

Work four of the seven following Griffiths' problems

3.4

3.10

3.14

3.27

3.30

## Additional Problems

All additional problems use the following two matrices where  $\hat{H}$  is the hamiltonian and  $\hat{A}$  is a second matrix associated with physical quantity  $a$ :

$$\hat{H} = \hbar\omega \begin{pmatrix} 1 & i & 0 \\ -i & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$\hat{A} = \hbar \begin{pmatrix} 1 & 0 & i \\ 0 & 1 & 0 \\ -i & 0 & 1 \end{pmatrix}$$

A1 Are either or both  $\hat{A}$  and  $\hat{H}$  Hermitian? Why? What does this imply?

A2 Calculate the uncertain relation,  $\sigma_E^2 \sigma_a^2 > ?$  for a system in the energy ground state.

A3 A system is prepared in the energy ground state. What values of  $a$  could be observed with what probability?

A4 A system is prepared in a linear combination of the energy ground state and the first excited state:

$$|\psi\rangle = \frac{1}{\sqrt{2}}|\phi_0\rangle + \frac{1}{\sqrt{2}}|\phi_1\rangle$$

Calculate the expectation value of  $a$ ,  $\langle a \rangle$ , as a function of time directly by calculating  $\langle \psi(t) |$  and then calculating the expectation value.

A5 A measurement is performed that finds the system is in a state with the lowest value of  $a$ . Calculate the expectation value of the energy for this state.

Each 25 pts

3.4 If  $\mathcal{I}(\hat{A}|b\rangle, |a\rangle) = \mathcal{I}(|b\rangle, \hat{A}|a\rangle)$

and  $\mathcal{I}(\hat{B}|b\rangle, |a\rangle) = \mathcal{I}(|b\rangle, \hat{B}|a\rangle)$

then  $\mathcal{I}(|b\rangle, (\hat{A} + \hat{B})|a\rangle) = \mathcal{I}(|b\rangle, \hat{A}|a\rangle)$   
 $+ \mathcal{I}(|b\rangle, \hat{B}|a\rangle)$

$$= \mathcal{I}(\hat{A}|b\rangle, |a\rangle) + \mathcal{I}(\hat{B}|b\rangle, |a\rangle)$$

$$= \mathcal{I}((\hat{A} + \hat{B})|b\rangle, |a\rangle)$$

$$\Rightarrow (\hat{A} + \hat{B})^\dagger = (\hat{A} + \hat{B})$$

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(b)  $(\alpha \hat{Q})^\dagger = \alpha^* \hat{Q}^\dagger = \alpha^* \hat{Q}$

$\Rightarrow \alpha$  real for Hermitian

(c)  $(\hat{A}\hat{B})^\dagger = \hat{B}^\dagger \hat{A}^\dagger = \hat{B}\hat{A}$

Hermitian if they commute

(d) The inner product for these operators

$$\text{is } \langle f | g \rangle = \int f^* g dx$$

Position  $\langle f | x g \rangle = \int f^* x g dx$

$$= \int x f^* g dx$$

$$= \langle x f | g \rangle \equiv \mathcal{I}(x f, g)$$

Hamiltonian

Since  $V(x)$  is just a function of  $x$ , the above reasoning shows it is Hermitian.

$$\mathcal{I}\left(f, -\frac{\hbar^2}{2m} \frac{d^2 g}{dx^2}\right) = -\frac{\hbar^2}{2m} \int_{-\infty}^{\infty} f^* \frac{d^2 g}{dx^2} dx$$

$$\mathcal{I}\left(f, \frac{d^2 g}{dx^2}\right) = \int_{-\infty}^{\infty} f^* \frac{d^2 g}{dx^2} dx$$

## Integrate by parts

$$\begin{aligned}\int f^* \frac{d^2 g}{dx^2} dx &= f^* \frac{dg}{dx} \Big|_{-\infty}^{\infty} - \int_{-\infty}^{\infty} \frac{df^*}{dx} \frac{dg}{dx} dx \\ &= f^* \frac{dg}{dx} \Big|_{-\infty}^{\infty} - g \frac{df^*}{dx} \Big|_{-\infty}^{\infty} + \int_{-\infty}^{\infty} g \frac{d^2 f^*}{dx^2} dx \\ &= \mathbb{I} \left( \frac{d^2 f}{dx^2}, g \right)\end{aligned}$$

We may assume surface terms vanish for square-integrable functions.

3.10

$$\phi_1 = \sqrt{\frac{2}{a}} \sin \frac{\pi x}{a}$$

$$\hat{p}_x = \frac{\hbar}{i} \frac{\partial}{\partial x}$$

$$\hat{p}_x \phi_1 = \left( \frac{\hbar}{i} \right) \sqrt{\frac{2}{a}} \frac{\pi}{a} \cos \frac{\pi x}{a}$$

$$\neq p \phi_1$$

Why? The ground state is formed by a left and right traveling wave.

3.27

(a) State collapses to  $\psi_1$

(b)  $\psi_1 = \frac{3}{5}\phi_1 + \frac{4}{5}\phi_2$  (Normalized)

Outcomes

$b_1$   $P(b_1) = \left(\frac{3}{5}\right)^2 = \frac{9}{25}$

$b_2$   $P(b_2) = \left(\frac{4}{5}\right)^2 = \frac{16}{25}$

(c) If  $b_1$  was observed, then state <sup>is</sup>  $\phi_1$

~~$4\psi_1 + 3\psi_2 = \frac{29}{5}\phi_1$~~

~~$4\psi_1 + 3\psi_2 =$~~

~~$\phi_1 = \frac{5}{6}\psi_1 + \frac{1}{6}\psi_2$~~

$$3\psi_1 + 4\psi_2 = \frac{25}{5}\phi_1$$

$$\phi_1 = \frac{3}{5}\psi_1 + \frac{4}{5}\psi_2$$

So the outcomes are

$$a_1 \quad P(a_1) = \frac{9}{25}$$

$$a_2 \quad P(a_2) = \frac{16}{25}$$

If  $b_2$  were observed, the state becomes  $\phi_2$

$$4\psi_1 - 3\psi_2 = \frac{25}{5}\phi_2$$

$$\phi_2 = \frac{4}{5}\psi_1 - \frac{3}{5}\psi_2$$

Outcomes

$$a_1 : P(a_1) = \frac{16}{25}$$

$$a_2 : P(a_2) = \frac{9}{25}$$

Total probability of observing  $a_1$  after  
SOME measurement of  $b$

$$P = \underbrace{\frac{9}{25}}_{b_1} \cdot \frac{9}{25} + \underbrace{\frac{16}{25}}_{b_2} \cdot \frac{16}{25}$$

$$= 0.54$$



3.30

(a) Normalize  $\psi = \frac{A}{x^2 + a^2}$

$$1 = \int_{-\infty}^{\infty} \frac{A^2}{(x^2 + a^2)^2} dx = A^2 \frac{\pi}{2a^3}$$

$$A = \sqrt{\frac{2a^3}{\pi}}$$

(b)  $\langle x \rangle = 0$  (~~odd~~ even function / odd range)

$$\langle x^2 \rangle = \int_{-\infty}^{\infty} x^2 \psi^* \psi dx = \int_{-\infty}^{\infty} \frac{A^2 x^2 dx}{(x^2 + a^2)^2}$$

$$= a^2$$

~~$$\sigma_x = \sqrt{\langle x^2 \rangle - \langle x \rangle^2} = a$$~~

$$\sigma_x = \sqrt{\langle x^2 \rangle - \langle x \rangle^2} = a$$

$$(c) \quad \overline{\Phi}(p, 0) = \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{\infty} \frac{A}{x^2 + a^2} e^{-ipx/\hbar} dx$$

$$= \frac{A}{\sqrt{2\pi\hbar}} \int_{-\infty}^{\infty} \frac{\cos px/\hbar - i \sin px/\hbar}{x^2 + a^2} dx$$

The sin piece is zero because of odd function / even range.

$$\overline{\Phi}(p, 0) = \frac{2A}{\sqrt{2\pi\hbar}} \int_0^{\infty} \frac{\cos px/\hbar}{x^2 + a^2} dx$$

=

## Schaum's

$$\int_0^{\infty} \frac{\cos mx}{x^2 + a^2} dx = \frac{\pi}{2a} e^{-ma}$$

$$\Phi(p, 0) = \frac{ZA}{\sqrt{2\pi\hbar}} \int_0^{\infty} \frac{\cos px/\hbar}{x^2 + a^2} dx$$

$$= \frac{ZA}{\sqrt{2\pi\hbar}} \frac{\pi}{2a} e^{-pa/\hbar}$$

$$= \frac{Z}{\sqrt{2\pi\hbar}} \frac{\pi}{2a} \sqrt{\frac{2a^3}{\pi}} e^{-pa/\hbar}$$

$$= \sqrt{\frac{a}{\hbar}} e^{-pa/\hbar}$$

Or if  $p < 0$ ,

$$\Phi(p, 0) = \sqrt{\frac{a}{\hbar}} e^{-|p|a/\hbar}$$

Maple gives

$$+ \sqrt{\frac{a}{\hbar}} \left( \cosh \frac{p_0}{\hbar} - \sinh \frac{p_0}{\hbar} \right)$$

$$= \sqrt{\frac{a}{\hbar}} e^{-p_0/\hbar}$$

Clearly not correct for  $p \rightarrow -\infty$

$$\Phi(p, 0) = \sqrt{\frac{a}{\hbar}} e^{-|p|a/\hbar}$$

Check normalization

$$1 = \frac{a}{\hbar} \int_{-\infty}^{\infty} e^{-2|p|a/\hbar} dp$$

$$= \frac{2a}{\hbar} \int_0^{\infty} e^{-2pa/\hbar} dp$$

$$= 1$$

(d)  $\langle p \rangle = 0$  odd function even range

$$\langle p^2 \rangle = \int_{-\infty}^{\infty} p^2 \Phi^* \Phi dp$$

$$= \frac{2a}{\hbar} \int_0^{\infty} p^2 e^{-2pa/\hbar} dp$$

$$= \frac{2a}{\hbar} \cdot \left( -\frac{1}{4} \right) \left( -\frac{\hbar^3}{a^3} \right)$$

$$= \frac{\hbar^2}{2a^2}$$

$$\sigma_p = \sqrt{\langle p^2 \rangle - \langle p \rangle^2} = \frac{1}{\sqrt{2}} \frac{\hbar}{a}$$

(e) Uncertainty

$$\sigma_x \sigma_p = a \cdot \frac{1}{\sqrt{2}} \frac{\hbar}{a} = \frac{1}{\sqrt{2}} \hbar = 0.707 \hbar > 0.5 \hbar$$

✓

$$\begin{aligned}
&> \text{int}\left(\frac{1}{(x^2 + a^2)^2}, x = -\text{infinity}..\text{infinity}\right); \\
&\quad \left\{ \begin{array}{ll} -\frac{1}{2} \frac{\pi}{a^3} & a < 0 \\ \infty & a = 0 \\ \frac{1}{2} \frac{\pi}{a^3} & 0 < a \end{array} \right. \quad (1)
\end{aligned}$$

$$\begin{aligned}
&> A := \left(\frac{2 \cdot a^3}{Pi}\right)^{\left(\frac{1}{2}\right)}; \\
&\quad A := \sqrt{2} \sqrt{\frac{a^3}{\pi}} \quad (2)
\end{aligned}$$

> assume(a > 0);

$$\begin{aligned}
&> \text{int}\left(\frac{A^2 \cdot x^2}{(x^2 + a^2)^2}, x = -\text{infinity}..\infty\right); \\
&\quad a^2 \quad (3)
\end{aligned}$$

$$\begin{aligned}
&> \frac{2 \cdot A}{(2 \cdot Pi \cdot h)^{\left(\frac{1}{2}\right)}} \cdot \text{int}\left(\frac{\cos\left(\frac{p \cdot x}{h}\right)}{(x^2 + a^2)}, x = 0..\text{infinity}\right); \\
&\quad \frac{1}{\sqrt{\pi h a^{\sim}}} \left( \sqrt{\frac{a^{\sim 3}}{\pi}} \left( \text{Icosh}\left(\frac{a^{\sim} - p}{h}\right) \text{Ci}\left(-\frac{1 a^{\sim} - p}{h}\right) - \text{Icosh}\left(\frac{a^{\sim} - p}{h}\right) \text{Ci}\left(\frac{1 a^{\sim} - p}{h}\right) \right. \right. \\
&\quad \left. \left. - \text{esgn}\left(\frac{p}{h}\right) \pi \sinh\left(\frac{a^{\sim} - p}{h}\right) \right) \right) \quad (4)
\end{aligned}$$

> simplify(%);

$$-\sqrt{a^{\sim}} \sqrt{\frac{1}{h}} \text{esgn}\left(\frac{p}{h}\right) \left( -\cosh\left(\frac{a^{\sim} - p}{h}\right) + \sinh\left(\frac{a^{\sim} - p}{h}\right) \right) \quad (5)$$

$$\begin{aligned}
&> \frac{2 \cdot a}{h} \cdot \text{int}\left(p^2 \cdot \exp\left(\frac{-2 \cdot p \cdot a}{h}\right), p = 0..\text{infinity}\right); \\
&\quad \frac{2 a^{\sim}}{h} \left( \lim_{p \rightarrow \infty} \left( \frac{-\frac{1}{4} h \left( e^{\left(\frac{-2 a^{\sim} - p}{h}\right)} h^2 + 2 e^{\left(\frac{-2 a^{\sim} - p}{h}\right)} p h a^{\sim} + 2 e^{\left(\frac{-2 a^{\sim} - p}{h}\right)} p^2 a^{\sim 2} - h^2 \right)}{a^{\sim 3}} \right) \right) \quad (6)
\end{aligned}$$

> simplify(%);

$$\frac{2 a^{\sim}}{h} \left( \lim_{p \rightarrow \infty} \left( \frac{-\frac{1}{4} h \left( e^{\left(\frac{-2 a^{\sim} - p}{h}\right)} h^2 + 2 e^{\left(\frac{-2 a^{\sim} - p}{h}\right)} p h a^{\sim} + 2 e^{\left(\frac{-2 a^{\sim} - p}{h}\right)} p^2 a^{\sim 2} - h^2 \right)}{a^{\sim 3}} \right) \right) \quad (7)$$

>

(A1) Both are Hermitian because

$$(\hat{H}^T)^* = \hat{H}^\dagger = \hat{H}$$

and  $(\hat{A}^T)^* = \hat{A}^\dagger = \hat{A}$

This implies the eigenvalues should be real  
and the eigenvectors should span, which they do.

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We will need the eigenvalues and eigenvectors of  
both matrices

Eigenvalues  $\hat{H}$   $0, \hbar\omega, 2\hbar\omega$

$$|0\hbar\omega\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} -i \\ 1 \\ 0 \end{pmatrix} \quad |1\hbar\omega\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

$$|2\hbar\omega\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} i \\ 1 \\ 0 \end{pmatrix}$$

Eigenvalues  $\hat{A}$

$0, \hbar, 2\hbar$

$$|0\hbar\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} -i \\ 0 \\ 1 \end{pmatrix}$$

$$|\hbar\rangle = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}$$

$$|2\hbar\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} i \\ 0 \\ 1 \end{pmatrix}$$



$\hat{A}$

eigenvectors  $\{\{1, 0, i\}, \{0, 1, 0\}, \{-i, 0, 1\}\}$

Assuming  $i$  is the imaginary unit | Use  $t$  as a variable instead

Input:

$$\text{Eigenvectors}\left[\begin{pmatrix} 1 & 0 & i \\ 0 & 1 & 0 \\ -i & 0 & 1 \end{pmatrix}\right]$$

Result:

$$v_1 \approx \{t, 0, 1\}$$

$$v_2 \approx \{0, 1, 0\}$$

$$v_3 \approx \{-t, 0, 1\}$$

Corresponding eigenvalues:

$$\lambda_1 = 2$$

$$\lambda_2 = 1$$

$$\lambda_3 = 0$$

### New to Wolfram|Alpha?

A few things to try:

enter any date (e.g. a birth date)  
june 23, 1988

enter any city (e.g. a home town)  
new york

enter any two stocks  
IBM Apple

enter any calculation  
\$250 + 15%

enter any math formula  
 $x^2 \sin(x)$

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$\hat{A}$

eigenvectors  $\{\{1, i, 0\}, \{-i, 1, 0\}, \{0, 0, 1\}\}$

Assuming  $i$  is the imaginary unit | Use  $i$  as a variable instead

Input:

$$\text{Eigenvectors}\left[\begin{pmatrix} 1 & i & 0 \\ -i & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}\right]$$

Result:

$$v_1 \approx \{i, 1, 0\}$$

$$v_2 \approx \{0, 0, 1\}$$

$$v_3 \approx \{-i, 1, 0\}$$

Corresponding eigenvalues.

$$\lambda_1 = 2$$

$$\lambda_2 = 1$$

$$\lambda_3 = 0$$

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In ground state

(A2)

$$\sigma_E^2 \sigma_A^2 \geq \left( \frac{1}{2i} \langle 0 | [\hat{H}, \hat{A}] | 0 \rangle \right)^2$$

Compute commutator

$$\hat{H} \hat{A} = \hbar^2 \omega \begin{pmatrix} 1 & i & 0 \\ -i & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 & i \\ 0 & 1 & 0 \\ -i & 0 & 1 \end{pmatrix}$$

$$= \hbar^2 \omega \begin{pmatrix} 1 & i & i \\ -i & 1 & 1 \\ -i & 0 & 1 \end{pmatrix}$$

$$\hat{A} \hat{H} = \hbar^2 \omega \begin{pmatrix} 1 & 0 & i \\ 0 & 1 & 0 \\ -i & 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & i & 0 \\ -i & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$= \hbar^2 \omega \begin{pmatrix} 1 & i & i \\ -i & 1 & 0 \\ -i & 1 & 1 \end{pmatrix}$$

$$[\hat{H}, \hat{A}] = \hat{H} \hat{A} - \hat{A} \hat{H} = \hbar^2 \omega \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & -1 & 0 \end{pmatrix}$$

$$\langle 0 | [H, A] | 0 \rangle$$

$$= \frac{1}{\sqrt{2}} (i \ 1 \ 0) \hbar^2 \omega \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & -1 & 0 \end{pmatrix} \begin{pmatrix} -i \\ 1 \\ 0 \end{pmatrix} \frac{1}{\sqrt{2}}$$

$$= \frac{\hbar^2 \omega}{2} (i \ 1 \ 0) \begin{pmatrix} 0 \\ 0 \\ -1 \end{pmatrix} = 0$$

For ground state,  $\sigma_E^2 \sigma_d^2 \geq 0$

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First excited state,  $\langle \hbar\omega | [H, A] | \hbar\omega \rangle$

$$= (0, 0 \ 1) \hbar^2 \omega \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & -1 & 0 \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

$$= \hbar^2 \omega (0 \ 0 \ 1) \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} = 0$$

$$\sigma_E^2 \sigma_d^2 \geq 0$$

A3

The energy ground state from (A1)

$$|0\hbar\omega\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} -i \\ 1 \\ 0 \end{pmatrix} = |\psi\rangle$$

The eigenvectors of  $\hat{A}$  also from (A1)

$$|0\hbar\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} -i \\ 0 \\ 1 \end{pmatrix} \quad |1\hbar\rangle = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}$$

$$|2\hbar\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} i \\ 0 \\ 1 \end{pmatrix}$$

The probability to observe  $0\hbar$  is

$$P(0\hbar) = |\langle 0\hbar | \psi \rangle|^2$$

$$\langle 0\hbar | \psi \rangle = \frac{1}{\sqrt{2}} (i \ 0 \ 1) \frac{1}{\sqrt{2}} \begin{pmatrix} -i \\ 1 \\ 0 \end{pmatrix} = \frac{1}{2} \cdot 1 = \frac{1}{2}$$

$$P(0\hbar) = \langle 0\hbar | \psi \rangle^* \cdot \langle 0\hbar | \psi \rangle = \frac{1}{4}$$

$$\langle \hbar | \psi \rangle = (0 \ 1 \ 0) \frac{1}{\sqrt{2}} \begin{pmatrix} -i \\ 1 \\ 0 \end{pmatrix} = \frac{1}{\sqrt{2}}$$

$$P(\hbar) = \langle \hbar | \psi \rangle^* \langle \hbar | \psi \rangle = \frac{1}{2}$$

$$\langle 2\hbar | \psi \rangle = \frac{1}{\sqrt{2}} (-i \ 0 \ 1) \frac{1}{\sqrt{2}} \begin{pmatrix} -i \\ 1 \\ 0 \end{pmatrix}$$

$$= \frac{1}{2} \cdot (-1) = -\frac{1}{2}$$

$$P(2\hbar) = \langle 2\hbar | \psi \rangle^2 \langle 2\hbar | \psi \rangle = \frac{1}{4}$$

Outcome (a)	Probability
$0\hbar = 0$	$\frac{1}{4}$
$\hbar$	$\frac{1}{2}$
$2\hbar$	$\frac{1}{4}$
Anything	1

(A4)

$$|\psi\rangle = \frac{1}{\sqrt{2}} |0\hbar\omega\rangle + \frac{1}{\sqrt{2}} |\hbar\omega\rangle$$

$$|\psi(t)\rangle = \hat{U} |\psi(0)\rangle$$

$$= \frac{1}{\sqrt{2}} |0\hbar\omega\rangle + \frac{1}{\sqrt{2}} e^{-i\omega t} |\hbar\omega\rangle$$

$$= \frac{1}{2} \begin{pmatrix} -i \\ 1 \\ 0 \end{pmatrix} + \frac{1}{\sqrt{2}} e^{-i\omega t} \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

$$= \begin{pmatrix} -i/2 \\ 1/2 \\ e^{-i\omega t}/\sqrt{2} \end{pmatrix}$$

Expectation Value

$$\langle 0 \rangle(t) = \langle \psi | \hat{A} | \psi \rangle$$

$$= \begin{pmatrix} \frac{i}{2} & \frac{1}{2} & \frac{e^{i\omega t}}{\sqrt{2}} \end{pmatrix} \begin{pmatrix} 1 & 0 & i \\ 0 & 1 & 0 \\ -i & 0 & 1 \end{pmatrix} \begin{pmatrix} -i/2 \\ 1/2 \\ \frac{e^{-i\omega t}}{\sqrt{2}} \end{pmatrix}$$

$$\langle a \rangle(t) = \hbar \left( \frac{i}{2} \quad \frac{1}{2} \quad \frac{e^{i\omega t}}{\sqrt{2}} \right) \begin{pmatrix} -\frac{1}{2} & + \frac{ie^{-i\omega t}}{\sqrt{2}} \\ & \frac{1}{2} \\ -\frac{1}{2} & + \frac{e^{-i\omega t}}{\sqrt{2}} \end{pmatrix}$$

$$= \hbar \left[ \frac{1}{4} - \frac{e^{-i\omega t}}{2\sqrt{2}} + \frac{1}{4} - \frac{e^{i\omega t}}{2\sqrt{2}} + \frac{1}{2} \right]$$

$$= \hbar \left[ 1 - \frac{1}{2\sqrt{2}} (e^{i\omega t} + e^{-i\omega t}) \right]$$

$$= \hbar \left[ 1 - \frac{1}{\sqrt{2}} \cos \omega t \right]$$



A5

$$|\psi\rangle = |0\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} -i \\ 0 \\ 1 \end{pmatrix}$$

Expectation Value of  $\hat{E}$

$$\langle E \rangle = \langle \psi | \hat{H} | \psi \rangle$$

$$= \frac{1}{\sqrt{2}} (i \ 0 \ 1) \hbar\omega \begin{pmatrix} 1 & i & 0 \\ -i & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \frac{1}{\sqrt{2}} \begin{pmatrix} -i \\ 0 \\ 1 \end{pmatrix}$$

$$= \frac{\hbar\omega}{2} (i \ 0 \ 1) \begin{pmatrix} -i \\ -1 \\ 1 \end{pmatrix}$$

$$= \frac{\hbar\omega}{2} (1 + 1) = \hbar\omega$$