

## SHO Again

The simple harmonic oscillator is

$$\hat{H} = \frac{\hat{P}^2}{2m} + \frac{1}{2} m \omega^2 \hat{X}$$

which has an energy spectrum

$$E_n = \left(n + \frac{1}{2}\right) \hbar \omega$$

and eigenvalue eqn

$$\hat{H} |\phi_n\rangle = E_n |\phi_n\rangle$$

We defined raising and lowering operators

$$\text{Raising } \hat{a}_+ \equiv \hat{a} = \frac{1}{\sqrt{2\hbar m \omega}} (-i\hat{P} + m\omega \hat{X})$$

$$\text{Lowering } \hat{a}_- \equiv \hat{a}^\dagger = \frac{1}{\sqrt{2\hbar m \omega}} (i\hat{P} + m\omega \hat{X})$$

②

The action of these operators move us between states,

$$\hat{a}_+ |\phi_n\rangle = \sqrt{n+1} |\phi_{n+1}\rangle$$

$$\hat{a}_- |\phi_n\rangle = \sqrt{n} |\phi_{n-1}\rangle$$

where  $\hat{a}_- |\phi_0\rangle = 0$ .

We can build any energy eigenstate by raising the ground state enough times

$$|\phi_n\rangle = \frac{1}{\sqrt{n!}} (\hat{a}_+)^n |\phi_0\rangle$$

What is  $|\phi_0\rangle$  is the  $\{ |x\rangle \}$  basis?

$$\hat{a}_- |\phi_0\rangle = 0$$

$$\frac{1}{\sqrt{2\hbar m\omega}} (i\hat{p} + m\omega\hat{x}) |\phi_0\rangle = 0$$

$$\frac{1}{\sqrt{2\hbar m\omega}} \left( i\frac{\hbar}{i} \frac{\partial}{\partial x} + m\omega x \right) \phi_0(x) = 0$$

$$\phi_0(x) = \langle x | \phi_0 \rangle$$

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$$\hbar \frac{d}{dx} \phi_0(x) + m\omega x \phi_0(x) = 0$$

$$\phi_0(x) = A e^{-\frac{m\omega}{2\hbar} x^2}$$

Normalize, which everyone has done for this function

$$\phi_0(x) = \left( \frac{m\omega}{\hbar} \right)^{1/4} e^{-\frac{m\omega}{2\hbar} x^2}$$

Therefore,

$$\phi_n = \frac{1}{\sqrt{n!}} (\hat{a}_+)^n |\phi_0\rangle$$

$$= \frac{1}{\sqrt{n!}} \left( \frac{1}{\sqrt{2\hbar m\omega}} \left[ -i \left( \hbar \frac{d}{dx} \right) + m\omega x \right] \right)^n \phi_0(x)$$

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With some work, or the use of Maple, this can be written

$$\phi_n(x) = \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} \frac{1}{\sqrt{2^n n!}} H_n(\xi) e^{-\xi^2/2}$$

$$\xi \equiv \sqrt{\frac{m\omega}{\hbar}} x$$

### Hermite Polynomials

$$H_0 = 1$$

$$H_1 = 2x$$

$$H_2 = 4x^2 - 2$$

$$H_3 = 8x^3 - 12x$$

So we may work in operator land or function space land

Expectation Value  $x$  in the ground state

$$\langle \phi_0 | \hat{X} | \phi_0 \rangle = \int_{-\infty}^{\infty} \phi_0^* x \phi_0 dx$$

$$= \sqrt{\frac{m\omega}{\pi\hbar}} \int_{-\infty}^{\infty} x e^{-\frac{m\omega}{\hbar} x^2} dx = 0$$

or  $\hat{X} = \sqrt{\frac{\hbar}{2m\omega}} (\hat{a}_+ + \hat{a}_-)$

$$\begin{aligned} \langle \phi_0 | \hat{X} | \phi_0 \rangle &= \sqrt{\frac{\hbar}{2m\omega}} \langle \phi_0 | \hat{a}_+ + \hat{a}_- | \phi_0 \rangle \\ &= \sqrt{\frac{\hbar}{2m\omega}} \left[ \langle \phi_0 | \hat{a}_+ | \phi_0 \rangle + \langle \phi_0 | \hat{a}_- | \phi_0 \rangle \right] \\ &\qquad\qquad\qquad \begin{matrix} = & & = \\ 0 & & 0 \\ \langle \phi_0 | \phi_1 \rangle & & \end{matrix} \\ &= 0 \end{aligned}$$

What about  $\langle \hat{X}^2 \rangle$ ?

$$\begin{aligned} \langle \hat{X}^2 \rangle &= \int_{-\infty}^{\infty} x^2 \phi_0^*(x) \phi_0(x) dx \\ &= \sqrt{\frac{m\omega}{\hbar}} \int_{-\infty}^{\infty} x^2 e^{-\frac{m\omega}{\hbar} x^2} dx \\ &= \langle \phi_0 | \hat{X}^2 | \phi_0 \rangle \\ &= \frac{\hbar}{2m\omega} \langle \phi_0 | (\hat{a}_+ + \hat{a}_-)^2 | \phi_0 \rangle \end{aligned}$$

(6)

$$\begin{aligned}
\langle \hat{x}^2 \rangle &= \frac{\hbar}{2m\omega} \langle \phi_0 | \hat{a}_+^2 + \hat{a}_-^2 + \hat{a}_+ \hat{a}_- + \hat{a}_- \hat{a}_+ | \phi_0 \rangle \\
&= \frac{\hbar}{2m\omega} \left[ \langle \phi_0 | \hat{a}_+^2 | \phi_0 \rangle + \langle \phi_0 | \hat{a}_-^2 | \phi_0 \rangle \right. \\
&\quad \left. + \langle \phi_0 | \hat{a}_+ \hat{a}_- | \phi_0 \rangle + \langle \phi_0 | \hat{a}_- \hat{a}_+ | \phi_0 \rangle \right] \\
&= \frac{\hbar}{2m\omega} \langle \phi_0 | \hat{a}_- \hat{a}_+ | \phi_0 \rangle
\end{aligned}$$

because

$$\begin{aligned}
\langle \phi_0 | \hat{a}_+^2 | \phi_0 \rangle &= \langle \phi_0 | \hat{a}_+ (\hat{a}_+ | \phi_0 \rangle) \\
&= \langle \phi_0 | \hat{a}_+ | \phi_1 \rangle = \sqrt{2} \langle \phi_0 | \phi_2 \rangle = 0
\end{aligned}$$

$$\langle \phi_0 | \hat{a}_-^2 | \phi_0 \rangle = \langle \phi_0 | \hat{a}_- (\hat{a}_- | \phi_0 \rangle) = 0$$

$$\langle \phi_0 | \hat{a}_+ (\hat{a}_- | \phi_0 \rangle) = 0$$

$$\hat{a}_+ |\phi_0\rangle = \sqrt{0+1} |\phi_1\rangle = |\phi_1\rangle$$

$$\hat{a}_- |\phi_1\rangle = \sqrt{1} |\phi_0\rangle = |\phi_0\rangle$$

~~$\langle \phi_0 | \hat{a}_- \hat{a}_+ | \phi_0 \rangle$~~   $\langle \phi_0 | \hat{a}_- \hat{a}_+ | \phi_0 \rangle = \langle \phi_0 | \hat{a}_- | \phi_1 \rangle$   
 $= \langle \phi_0 | \phi_0 \rangle = 1$

$$\langle x^2 \rangle = \frac{\hbar}{2mw}$$

3.34 The ground state has energy  $\frac{1}{2}\hbar\omega$  and the first excited state  $\frac{3}{2}\hbar\omega$ . At  $t=0$ ,

$$|\psi\rangle = \frac{1}{\sqrt{2}} (|0\rangle + e^{i\theta}|1\rangle)$$

Since we know the probabilities are equal but we don't know the phase relationship.

$$\text{At time } t, \quad |\psi\rangle = \frac{1}{\sqrt{2}} \left( e^{-i\omega_0 t} |0\rangle + e^{-i\omega_0 t + i\theta} |1\rangle \right)$$

where I have introduced  $\omega_0 = \frac{1}{2}\omega$

$$|\psi\rangle = \frac{e^{-i\omega_0 t}}{\sqrt{2}} \left( |0\rangle + e^{-2i\omega_0 t + i\theta} |1\rangle \right)$$

From chapter 2

$$p = i\sqrt{\frac{\hbar m \omega}{2}} (a_+ - a_-)$$

2.69

$$a_+ |n\rangle = \sqrt{n+1} |n+1\rangle$$

$$a_- |n\rangle = \sqrt{n} |n-1\rangle$$

2.66



$$\text{Let } A = \frac{e^{-i\omega_0 t}}{\sqrt{2}} \quad \gamma = e^{-2i\omega_0 t + i\theta}$$

$$|\psi\rangle = A(|0\rangle + \gamma|1\rangle)$$

$$p|\psi\rangle = \left( A i \sqrt{\frac{\hbar m \omega}{2}} \right) (\alpha_+ |0\rangle + \gamma \alpha_+ |1\rangle - \gamma \alpha_- |1\rangle)$$

$$= \left( A i \sqrt{\frac{\hbar m \omega}{2}} \right) (|1\rangle + \gamma \sqrt{2} |2\rangle - \gamma |0\rangle)$$

$$\langle \psi | p | \psi \rangle = \left( A^* A i \sqrt{\frac{\hbar m \omega}{2}} \right) (\langle 0 | + \gamma^* \langle 1 |) (|1\rangle + \gamma \sqrt{2} |2\rangle - \gamma |0\rangle)$$

$$= \left( \frac{i}{2} \sqrt{\frac{\hbar m \omega}{2}} \right) (\gamma^* - \gamma)$$

$$\gamma^* - \gamma = +2i \sin(2\omega_0 t - \theta)$$

$$\langle \psi | p | \psi \rangle = - \sqrt{\frac{\hbar m \omega}{2}} \sin(2\omega_0 t - \theta)$$

$$= - \sqrt{\frac{\hbar m \omega}{2}} \sin(\omega t - \theta)$$

$\langle p \rangle$ 's maximum value is  $\sqrt{\frac{\hbar m \omega}{2}}$

If this is maximum at  $t=0$ ,  $\theta = +\frac{\pi}{2}$

The - sign out front means  $\langle p \rangle$  maximum when  
sin minimum.

$$|\psi\rangle = \frac{e^{-i\omega t/2}}{\sqrt{2}} (|0\rangle + i|1\rangle)$$

2.41

$$\phi_0 = \left( \frac{m\omega}{\pi\hbar} \right)^{1/4} e^{-\xi^2/2} = B e^{-\xi^2/2}$$

$$\phi_1 = \left( \frac{m\omega}{\pi\hbar} \right)^{1/4} \frac{2\xi}{\sqrt{2}} e^{-\xi^2/2} = B\sqrt{2}\xi e^{-\xi^2/2}$$

$$\phi_2 = \left( \frac{m\omega}{\pi\hbar} \right)^{1/4} \frac{4\xi^2 - 2}{2\sqrt{2}} e^{-\xi^2/2} \quad \text{from 2.85}$$

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$$= B(\sqrt{2}\xi^2 - \frac{1}{\sqrt{2}}) e^{-\xi^2/2}$$

$$\xi = \sqrt{\frac{m\omega}{\hbar}} x$$

$$B = \left( \frac{m\omega}{\pi\hbar} \right)^{1/4}$$

The given wave function is

$$\psi(x,0) = A(1 - 2\xi)^2 e^{-\xi^2/2}$$

$$= A(1 - 4\xi + 4\xi^2) e^{-\xi^2/2}$$

$$= c_0\phi_0 + c_1\phi_1 + c_2\phi_2$$

$$\Rightarrow 4A = c_2 B\sqrt{2}$$

Equate powers of  $\xi$

$$c_2 = 2\sqrt{2} \frac{A}{B}$$

$$-4A = B\sqrt{2}c_1$$

$$c_1 = -2\sqrt{2}\frac{A}{B}$$

$$\xi^0: A = c_0 B - c_2 \frac{B}{\sqrt{2}}$$

$$= c_0 B - 2\sqrt{2}\frac{A}{B} \left( \frac{B}{\sqrt{2}} \right)$$

$$= c_0 B - 2A$$

$$c_0 = \frac{3A}{B}$$

Normalize to find A

$$c_0^* c_0 + c_1^* c_1 + c_2^* c_2 = 1$$

$$\frac{9A^2}{B^2} + \frac{8A^2}{B^2} + \frac{8A^2}{B^2} = 1$$

$$A = \frac{B}{\sqrt{25}} = \frac{B}{5}$$

$$c_0 = \frac{3}{5} \quad c_1 = -\frac{2\sqrt{2}}{5} \quad c_2 = \frac{2\sqrt{2}}{5}$$

$$\psi = \frac{3}{5} \phi_0 - \frac{2\sqrt{2}}{5} \phi_1 + \frac{2\sqrt{2}}{5} \phi_2$$

$$\begin{aligned} \langle E \rangle &= c_0^* c_0 \left(\frac{1}{2} \hbar \omega\right) + c_1^* c_1 \left(\frac{3}{2} \hbar \omega\right) + c_2^* c_2 \left(\frac{5}{2} \hbar \omega\right) \\ &= \frac{9}{25} \left(\frac{1}{2} \hbar \omega\right) + \frac{8}{25} \left(\frac{3}{2} \hbar \omega\right) + \frac{8}{25} \left(\frac{5}{2} \hbar \omega\right) \\ &= \frac{\hbar \omega}{50} (9 + 24 + 40) \end{aligned}$$

$$\boxed{\langle E \rangle = \frac{73}{50} \hbar \omega}$$

$$\begin{aligned} (b) \quad \psi(x,t) &= \frac{3}{5} e^{-i\omega t/2} \phi_0 - \frac{2\sqrt{2}}{5} e^{-3i\omega t/2} \phi_1 \\ &\quad + \frac{2\sqrt{2}}{5} e^{-5i\omega t/2} \phi_2 \end{aligned}$$

~~The change affects only the  $\phi_1$  term changing its sign~~

~~301~~  
~~2~~

Factor out the overall phase

$$\psi(x, t) = e^{-i\omega t/2} \left[ \frac{3}{5} \phi_0 - \frac{2\sqrt{2}}{5} e^{-i\omega t} \phi_1 + \frac{2\sqrt{2}}{5} e^{-2i\omega t} \phi_2 \right]$$

We want to change the sign of the second term  
without affecting the third term

$$\omega T = \pi$$

$$T = \frac{\pi}{\omega}$$