

## Section - Rigid Bodies

Rigid Body is a system of particles whose relative positions are fixed.

Center of Mass  $\vec{r}_{cm} = (x_{cm}, y_{cm}, z_{cm})$

$$x_{cm} = \frac{1}{M} \sum_i x_i m_i$$

or

$$x_{cm} = \frac{1}{M} \int x \, dm$$

where  $dm = \rho \, dV$  or  $\sigma \, dA$  or  $\lambda \, dl$

where  $\rho \equiv \text{mass/volume}$

$\sigma \equiv \text{mass/area}$

$\lambda \equiv \text{mass/length}$ .

## Center of Mass of Composite Body - Suppose

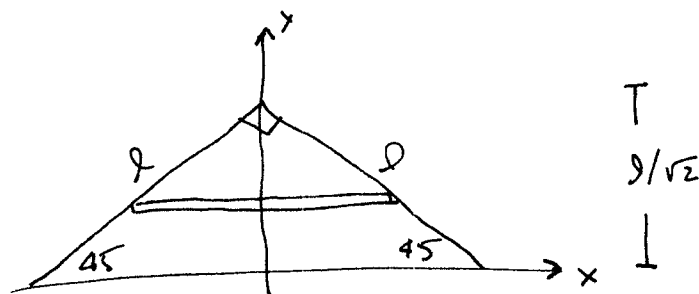
a body is made of separate parts with known centers of mass  $m_1, \vec{r}_{1cm}, m_2, \vec{r}_{2cm}$

then 
$$\vec{r}_{cm} = \frac{1}{M} \sum m_i \vec{r}_{i,cm}$$

$\Rightarrow$  Evident from definition of CM.

Symmetry - If a body has a line or plane of symmetry, the center of mass lies on that line or plane.

Example Center of Mass of isosceles right triangular lamina, mass density  $\sigma$ .



Total mass  $M = \frac{1}{2} \text{base} \times \text{height} = \frac{1}{2} l^2 \sigma$

By symmetry the center of mass lies on the  $y$ -axis,  $x_{cm} = 0$

$$y_{cm} = \frac{1}{M} \int y \, dm$$

The length of the strip drawn is  $z(l/\sqrt{2} - y)$

$$\cancel{y} \, dm = (\sigma)(z)(l/\sqrt{2} - y) \, dy$$

$$y_{cm} = \frac{1}{M} \int y \, dm = \frac{z\sigma}{M} \int_0^{l/\sqrt{2}} y(l/\sqrt{2} - y) \, dy$$

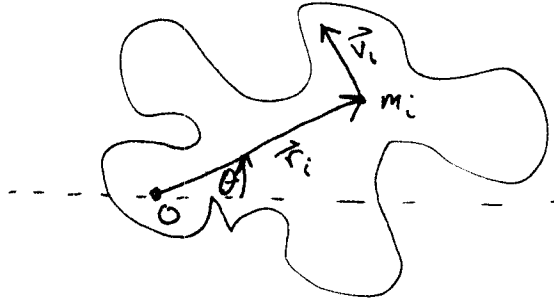
$$= \frac{z\sigma}{M} \left[ \frac{y^2 l}{2\sqrt{2}} - \frac{y^3}{3} \right]_0^{l/\sqrt{2}}$$

$$= \frac{z\sigma}{M} \left[ \frac{l^3}{4\sqrt{2}} - \frac{l^3}{6\sqrt{2}} \right]$$

$$= \frac{z\sigma l^3}{M} \left[ \frac{1}{12\sqrt{2}} \right]$$

$$y_{cm} = \frac{l}{3\sqrt{2}}$$

## Section - Rotation about Fixed Axis



If the body is rigid and is constrained to rotate about fixed axis  $O$  ( $z$  axis). All the particles move in circles with angular velocity  $\omega$  and speed  $v_i = r_i \omega$

$$\vec{r}_i = r_i (\cos \theta, \sin \theta, 0)$$

$$\vec{v}_i = r_i (-\sin \theta \dot{\theta}, \cos \theta \dot{\theta}, 0)$$

$$= \omega r_i (-\sin \theta, \cos \theta, 0)$$

Observation gives, where  $\vec{\omega} = \omega \hat{k}$

$$\boxed{\vec{v}_i = \vec{\omega} \times \vec{r}_i}$$

## Kinetic Energy (Pure Rotation)

$$T_{\text{rot}} = \frac{1}{2} \sum_i m_i v_i^2 = \frac{\omega^2}{2} \sum_i m_i r_i^2$$

Define Moment of Inertia about axis of rotation  $z$

$$I_z = \frac{1}{2} \sum_i m_i r_i^2 = \frac{1}{2} \sum_i m_i (x_i^2 + y_i^2)$$

## Angular Momentum

$$\vec{L} = \sum_i \vec{L}_i = \sum_i m_i \vec{r}_i \times \vec{v}_i$$

$$\vec{L}_i = m_i \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ x_i & y_i & 0 \\ \dot{x}_i & \dot{y}_i & 0 \end{vmatrix}$$

$$= m_i (x_i \dot{y}_i - \dot{y}_i x_i) \hat{k}$$

$$\vec{v}_i = r_i (\sin \theta, \cos \theta) \omega$$

$$\dot{x}_i = -r_i \omega \sin \theta \quad \dot{y}_i = r_i \omega \cos \theta$$

$$\vec{L}_i = m_i r_i \omega (x_i r_i \cos\theta + y_i r_i \sin\theta)$$
$$= m_i \omega (x_i^2 + y_i^2)$$

$$\vec{L} = \omega \sum_i m_i (x_i^2 + y_i^2) = I_z \omega$$

$$\underline{\text{Torque}} \quad \frac{dN_z}{dt} = \frac{dI_z \omega}{dt} = I_z \dot{\omega}$$

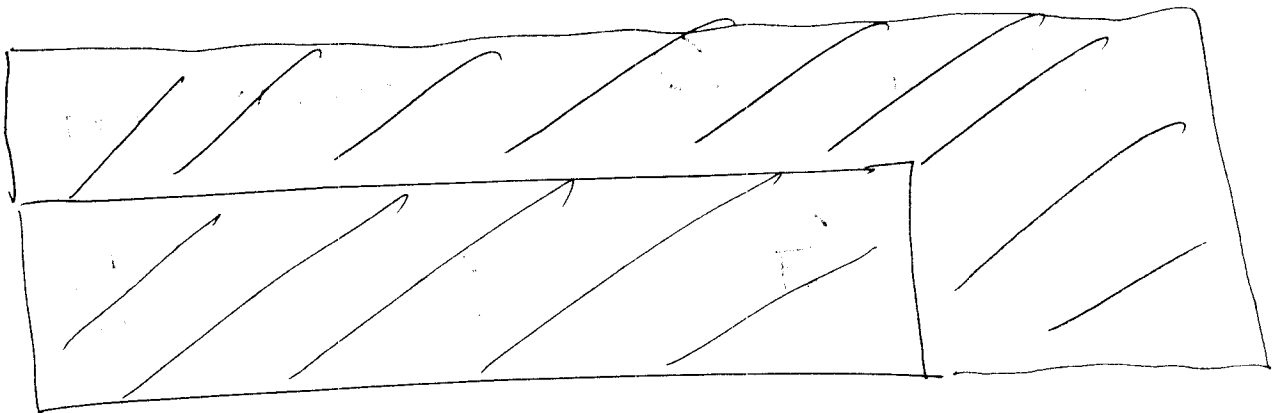
Lecture 3/28/2003

# Section - Comparison of Rocket Methods

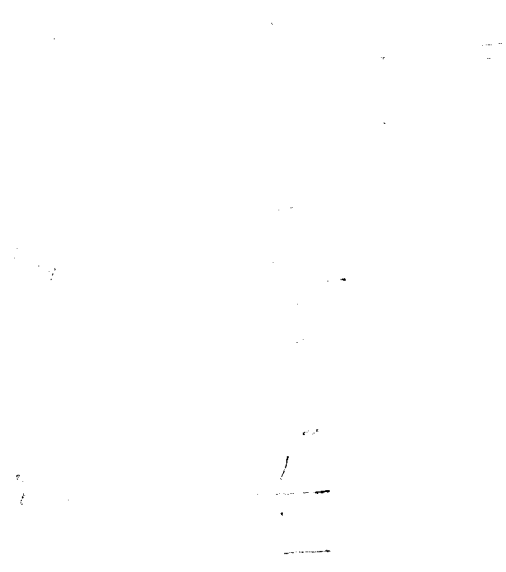
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Let mass per unit length be  $\lambda$

Lifted Mass $\lambda m$
Total Mass $\lambda L$



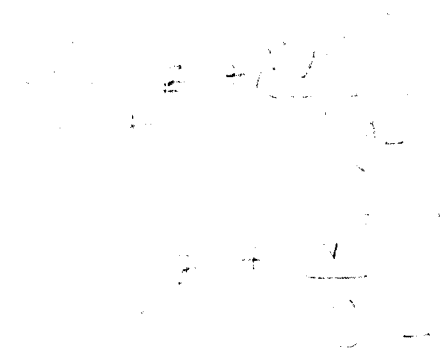




Force of gravity and normal force cancel for rope lying on table.

$$Mg = M_{\text{cm}}g$$

M:



1.  $\frac{d}{dt} (m \mathbf{v}) = \mathbf{F}$

2.  $\frac{d}{dt} (m \mathbf{v}) = \mathbf{F}$   
 $\frac{d}{dt} (m) \mathbf{v} + m \frac{d}{dt} (\mathbf{v}) = \mathbf{F}$

3.  $\frac{d}{dt} (m \mathbf{v}) = \mathbf{F}$   
 $\frac{d}{dt} (m) \mathbf{v} + m \frac{d}{dt} (\mathbf{v}) = \mathbf{F}$

4.  $\frac{d}{dt} (m \mathbf{v}) = \mathbf{F}$

$$\frac{d}{dt} (m \mathbf{v}) = \mathbf{F}$$

\* Adding Mass

$$\frac{d}{dt} (m \mathbf{v}) = \mathbf{F}$$

\* Added mass moves with negative velocity

$$\frac{d}{dt} (m \mathbf{v}) = \mathbf{F}$$

$$\frac{d}{dt} (m \mathbf{v}) = \mathbf{F}$$

## Section - Rigid Bodies

Rigid Body - A system of particles whose relative positions are fixed.

This means in the coordinate system where the body is stationary,  $\vec{r}_i - \vec{r}_j$  does not depend on time.

Center of Mass

$$\vec{r}_{cm} = (x_{cm}, y_{cm}, z_{cm})$$

$$x_{cm} = \frac{1}{M} \sum x_i m_i$$

- or -

$$x_{cm} = \frac{1}{M} \int x \, dm$$

$dm$  is an element of mass

$$dm = \underline{\rho \, dV} \quad \text{or} \quad \underline{\sigma \, dA} \quad \text{or} \quad \underline{\lambda \, dl}$$

where  $\rho$  is the volume mass density  
 $\sigma$  is the surface mass density  
 $\lambda$  is the linear mass density

Center of Mass of Composite Body - Suppose

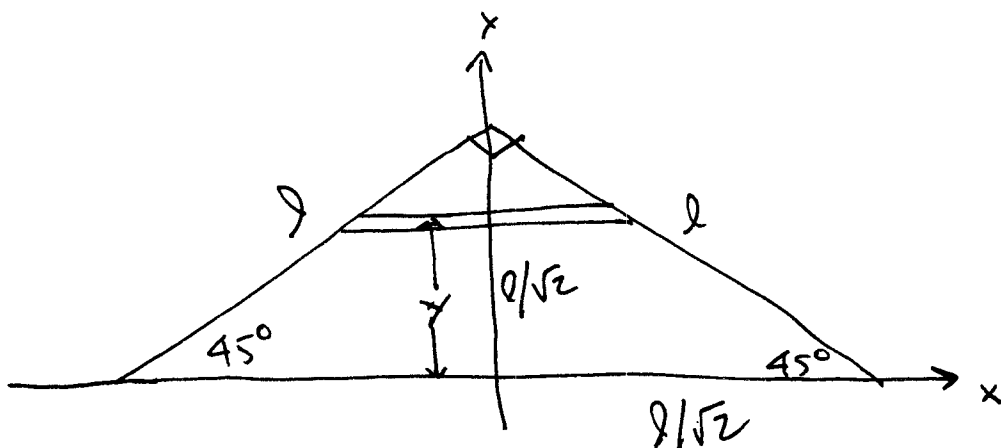
a body is made of  $N$  parts whose masses are  $m_1, \vec{r}_1^{cm}$  and  $m_2, \vec{r}_2^{cm} \dots$

Then the Center of Mass of the composite body is

$$\vec{r}_{cm} = \frac{1}{m_1 + m_2 \dots} \left[ m_1 \vec{r}_1^{cm} + m_2 \vec{r}_2^{cm} \dots \right]$$

Symmetry - If a body has a line or plane of symmetry, the CM lies on that line/plane.

Example - Find CM of isosceles right planar triangle with uniform surface mass density  $\sigma$ .



Total Mass  $M = \frac{1}{2} \text{ base} \cdot \text{height} \cdot \sigma$   
 $= \frac{1}{2} \sigma l^2$

Symmetry  $xz$ -plane,  $yz$ -plane (or plane)  
 $x_{cm} = 0$   $y_{cm} = 0$

Definition

$$M x_{cm} = \int x \, dm$$

$$dm = (\text{length of strip}) \, dy \, \sigma$$

$$= 2\sigma \left( \frac{l}{\sqrt{2}} - y \right) dy$$

$$M x_{cm} = 2\sigma \int_0^{l/\sqrt{2}} \left( \frac{l}{\sqrt{2}} - y \right) y \, dy$$

$$= 2\sigma \left[ \frac{l}{\sqrt{2}} \frac{y^2}{2} - \frac{y^3}{3} \right]_0^{l/\sqrt{2}}$$

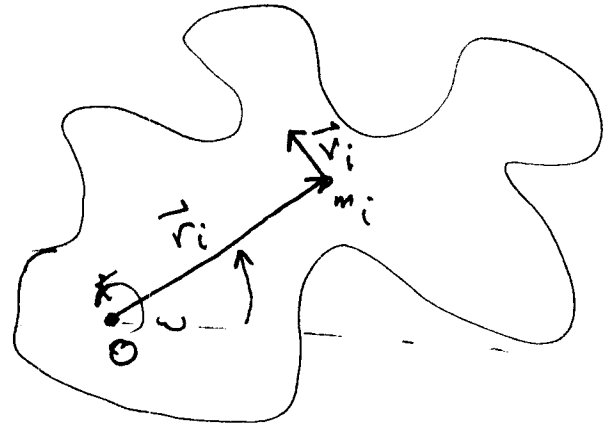
$$= 2\sigma \left[ \frac{l^3}{4\sqrt{2}} - \frac{l^3}{6\sqrt{2}} \right]$$

$$= \frac{\sigma l^3}{6\sqrt{2}} = \frac{\sigma l^3}{6\sqrt{2}}$$

$$\gamma_{cm} = \frac{\frac{\sigma \rho^3}{\sqrt{2}}}{\frac{1}{2} \sigma \rho^2} = \frac{\rho}{\sqrt{2}} = \frac{\rho}{3\sqrt{2}}$$

# Section - Rotation of Rigid Body about Fixed Axis

Motion of a plane object



If a body is rigid and constrained to rotate about the point O with angular velocity  $\vec{\omega}$ .

Velocity of any point  $v_i = r_i \omega$   $\omega = \dot{\theta}$

$$\vec{r}_i = r_i (\cos \theta, \sin \theta, 0)$$

$$\begin{aligned} \vec{v}_i &= \dot{\vec{r}}_i = r_i (-\sin \theta \dot{\theta}, \cos \theta \dot{\theta}, 0) \\ &= r_i \omega (-\sin \theta, \cos \theta, 0) = r_i \omega \hat{e}_\theta \end{aligned}$$

This can be written

$$\vec{v}_i = \vec{\omega} \times \vec{r}_i$$

Check Magnitude

$$|\vec{v}_i| = |\vec{\omega} \times \vec{r}_i| = \omega r_i$$

Check Direction

Use RHR, direction correct.

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Kinetic Energy (Pure Rotation) -

$$T_{\text{rot}} = \frac{1}{2} \sum_i m_i v_i^2$$

$$|\vec{v}_i| = \omega r_i$$

$$T_{\text{rot}} = \frac{\omega^2}{2} \sum_i m_i r_i^2 = \frac{1}{2} I \omega^2$$

Moment of Inertia (About O) -

$$I = \sum_i m_i r_i^2 = \sum_i m_i (x_i^2 + y_i^2)$$



# Angular Momentum ( $\vec{L}$ )

$$\vec{L} = \sum_i \vec{L}_i = \sum_i m_i \vec{r}_i \times \vec{v}_i$$

$$\vec{L}_i = m_i \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ x_i & y_i & 0 \\ \dot{x}_i & \dot{y}_i & 0 \end{vmatrix}$$

$$= m_i (x_i \dot{y}_i - y_i \dot{x}_i) \hat{k}$$

$$\dot{r}_i = \vec{v}_i = r_i \omega (-\sin \theta, \cos \theta, 0)$$

$$\dot{x}_i = -r_i \omega \sin \theta \quad \dot{y}_i = r_i \omega \cos \theta$$

$$\vec{L}_i = m_i (x_i r_i \omega \cos \theta + y_i r_i \omega \sin \theta) \hat{k}$$

$$x_i = r_i \cos \theta \quad y_i = r_i \sin \theta$$

$$\vec{L}_i = m_i \omega (x_i^2 + y_i^2) \hat{k}$$

$$\vec{L} = \sum \vec{L}_i = \omega \sum m_i (x_i^2 + y_i^2) \hat{k} = I \omega \hat{k}$$

# EOM

~~$\vec{N} = \frac{d\vec{L}}{dt}$~~

$$\vec{N} = \frac{d\vec{L}}{dt}$$

z-component

$$N_z = \frac{d I \omega}{dt} = I \dot{\omega}$$

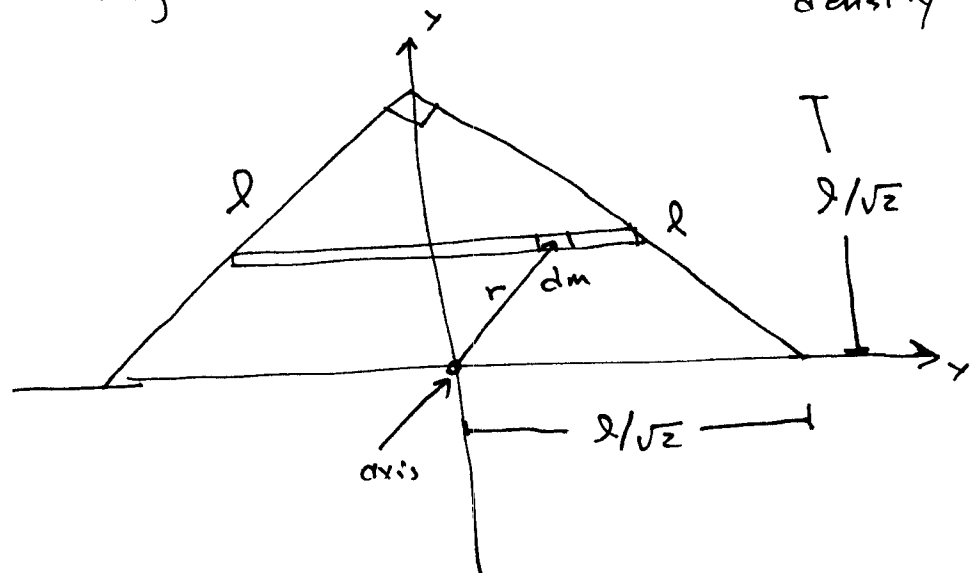
## Section - Moment of Inertia

Moment of Inertia - Moment of Inertia about any axis

$$I = \int dm r^2$$

where  $r^2$  is the perpendicular distance from the axis.

Example Moment of Inertia of Isosceles Right triangle about  $z$  axis. Uniform mass density  $\sigma$ .



Cut triangle into strips -

$$I = \int r^2 dm$$

$$dm = \sigma dx dy$$

$$r^2 = x^2 + y^2$$

$$= \sigma \int_0^{l/\sqrt{2}} dy \int_{-(l/\sqrt{2}-y)}^{l/\sqrt{2}-y} dx (x^2 + y^2)$$

$$= 2\sigma \int_0^{l/\sqrt{2}} dy \left[ \frac{x^3}{3} + xy^2 \right]_0^{l/\sqrt{2}-y}$$

$$= 2\sigma \int_0^{l/\sqrt{2}} dy \left[ \frac{1}{3} (l/\sqrt{2} - y)^3 + \frac{l}{\sqrt{2}} y^2 - y^3 \right]$$

Ask if  
continue.

$$= 2\sigma \int_0^{l/\sqrt{2}} \left[ \frac{1}{3} \left( \frac{l^3}{2\sqrt{2}} - \frac{3l^2}{2} y + \frac{3l}{\sqrt{2}} y^2 - y^3 \right) + \frac{l}{\sqrt{2}} y^2 - y^3 \right] dy$$

$$= 2\sigma \int_0^{l/\sqrt{2}} \left[ \frac{l^3}{6\sqrt{2}} - \frac{l^2}{2} y + \frac{2l}{\sqrt{2}} y^2 - \frac{4}{3} y^3 \right] dy$$

$$= 2\sigma \left( \frac{l^3}{6\sqrt{2}} y - \frac{l^2 y^2}{4} + \frac{2l y^3}{3\sqrt{2}} - \frac{1}{3} y^4 \right)_0^{l/\sqrt{2}}$$

Perpendicular Axis Thm - For a thin lamina  
in the  $x$ - $y$  plane.

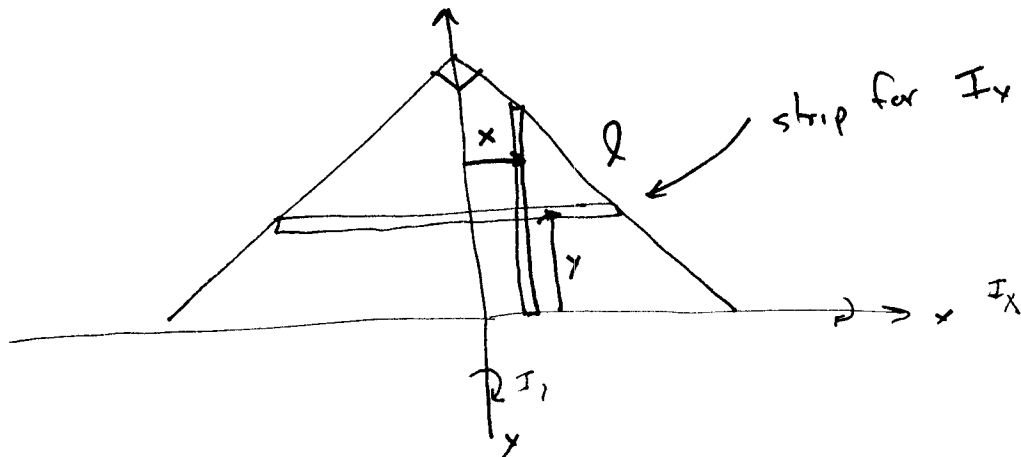
$$I_z = I_x + I_y$$

$$= \underbrace{\sum m_i (x_i^2 + y_i^2)}_{I_z} = \underbrace{\sum m_i x_i^2}_{I_y} + \underbrace{\sum m_i y_i^2}_{I_x}$$

The moment of inertia about an axis perpendicular to the lamina is equal to the sum of the moments of two <sup>mutually</sup> perpendicular axis through the axis.

Example Back to Right Isosceles Triangle

$$I_z = I_x + I_y$$



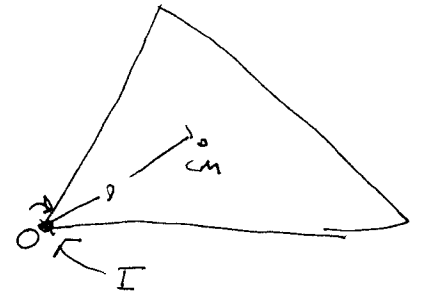
$$I_x = \int dm y^2 = 2 \int_0^{l/\sqrt{2}} \sigma (l/\sqrt{2} - y) y^2 dy$$

$$I_y = \int x^2 dm = \sigma \int_{-l/\sqrt{2}}^{l/\sqrt{2}} dx x^2 (l/\sqrt{2} - x)$$

## Parallel Axis Theorem

The moment of inertia about an axis

$O$  is the sum of the moment of inertia about a parallel axis through the center of mass plus the moment of inertia of the body, treated as a point mass at the center of mass about axis  $O$ .



If  $l$  is the distance from  $O$  to  $cm$ ,

$$I = I_{cm} + ml^2$$

Radius of Gyration ( $k$ ) Distance a point mass must be placed from the axis of rotation for the point mass to have the same moment of inertia as the body - Since

$$I_{\text{point}} = mk^2$$

$$k = \sqrt{\frac{I}{m}}$$